# Hamiltonian diffeomorphisms of toric manifolds and flag manifolds 

Andrés Viña<br>Departamento de Física, Universidad de Oviedo, Avda Calvo Sotelo, 33007 Oviedo, Spain<br>Received 11 January 2006; received in revised form 3 July 2006; accepted 16 July 2006<br>Available online 17 August 2006


#### Abstract

If $s$ and $n$ are integers relatively prime and $\operatorname{Ham}\left(G_{s}\left(\mathbb{C}^{n}\right)\right)$ is the group of Hamiltonian symplectomorphisms of the Grassmannian manifold $G_{s}\left(\mathbb{C}^{n}\right)$, we prove that $\sharp \pi_{1}\left(\operatorname{Ham}\left(G_{s}\left(\mathbb{C}^{n}\right)\right)\right) \geq n$.

We prove that $\pi_{1}(\operatorname{Ham}(M))$ contains an infinite cyclic subgroup, when $M$ is the one point blow up of $\mathbb{C} P^{3}$. We give a sufficient condition for the group $\pi_{1}(\operatorname{Ham}(M))$ to contain a subgroup isomorphic to $\mathbb{Z}^{p}$, when $M$ is a general toric manifold.


(c) 2006 Elsevier B.V. All rights reserved.

MSC: 53D05; 57S05
Keywords: Hamiltonian diffeomorphisms; Symplectic fibrations; Toric manifolds

## 1. Introduction

Let $(M, \omega)$ be a closed symplectic $2 n$-manifold. $\operatorname{By} \operatorname{Ham}(M, \omega)$ is denoted the group of Hamiltonian symplectomorphisms of $(M, \omega)$ [21,25]. The homotopy type of $\operatorname{Ham}(M, \omega)$ is only completely known in a few particular cases [20,25]. When $M$ is a surface, $\operatorname{Diff}_{0}(M)$ (the connected component of the identity map in the diffeomorphism group of $M$ ) is homotopy equivalent to the symplectomorphism group of $M$, hence the topology of the groups Ham $(M)$ in dimension 2 can be deduced from the description of the diffeomorphism groups of surfaces given in [6] (see [25]). On the other hand, positivity of the intersections of $J$-holomorphic spheres in 4-manifolds played a crucial role in the proof of results about the homotopy type of $\operatorname{Ham}(M)$ when $M$ is a ruled surface (see $[9,1$, 2]). But these arguments which work in dimension 2 or dimension 4 cannot be generalized to higher dimensions.

Using properties of the symplectic action on quantizable manifolds, in [27] we gave a lower bound for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{Q}))$, when $\mathcal{Q}$ is a quantizable coadjoint orbit of a semisimple Lie group $G$ and this orbit satisfies some technical hypotheses. By quite different methods McDuff and Tolman have proved the following result: if $\mathcal{O}$ is a coadjoint orbit of the semisimple group $G$, and the action of $G$ on $\mathcal{O}$ is effective, then the inclusion $G \rightarrow \operatorname{Ham}(\mathcal{O})$ induces an injection on $\pi_{1}$ [23]. This result answers a question posed in [29].

Here we give a new approach to determine lower bounds for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O}))$. We will consider curves in $G$ with initial point at $e$. Every family $\left\{g_{t} \mid t \in[0,1]\right\}$ of elements of $G$, with $g_{0}=e$ and $g_{1} \in Z(G)$ defines a loop $\psi$ in the group $\operatorname{Ham}(\mathcal{O})$. We will use the Maslov index of the linearized flow to deduce conditions under which the loops

[^0]$\psi$ and $\tilde{\psi}$, generated by the families $g_{t}$ and $\tilde{g}_{t}$ with different endpoints, are not homotopic. So our lower bounds for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O}))$ will be no bigger than $\sharp Z(G)$.

In particular we consider the group $S U(n+1)$ and an orbit $\mathcal{O}$ diffeomorphic to the Grassmannian $G_{s}\left(\mathbb{C}^{n+1}\right)$ of $s$-dimensional subspaces in $\mathbb{C}^{n+1}$, with $s$ and $n+1$ relatively prime. For each element of $Z(S U(n+1))$ we will construct a curve $g_{t}$ in $S U(n+1)$ such that the corresponding loops in $\operatorname{Ham}(\mathcal{O})$ relative to different elements of $Z(S U(n+1))$ have distinct Maslov indices. So these loops are homotopically inequivalent, and we have the following theorem.

Theorem 1. If $\mathcal{O}$ is a coadjoint orbit of $S U(n+1)$ diffeomorphic to the Grassmannian $G_{s}\left(\mathbb{C}^{n+1}\right)$ with $s$ and $n+1$ relatively prime, then

$$
\begin{equation*}
\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O})) \geq n+1 . \tag{1.1}
\end{equation*}
$$

In [27] we gave a lower bound for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{Q}))$, when $\mathcal{Q}$ is a quantizable coadjoint orbit. For the Grassmann manifolds to which the results of [27] are applicable, the bound given in [27] coincides with (1.1). However the results obtained here are more general, since now we do not assume that $\mathcal{O}$ is quantizable. We also give a lower bound for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O}))$ when $\mathcal{O}$ is a coadjoint orbit of $S U(n+1)$ diffeomorphic to a general flag manifold in $\mathbb{C}^{n+1}$.

On the other hand, a loop $\psi$ in the group $\operatorname{Ham}(M, \omega)$ determines a Hamiltonian fibration $E \xrightarrow{\pi} S^{2}$ with standard fibre $M$. On the total space $E$ we can consider the first Chern class $c_{1}(V T E)$ of the vertical tangent bundle of $E$. Moreover on $E$ is also defined the coupling class $c_{\psi} \in H^{2}(E, \mathbb{R})$ [11]. This class is determined by the following properties:
(i) $i_{q}^{*}\left(c_{\psi}\right)$ is the cohomology class of the symplectic structure on the fibre $\pi^{-1}(q)$, where $i_{q}$ is the inclusion of $\pi^{-1}(q)$ in $E$ and $q$ is an arbitrary point of $S^{2}$.
(ii) $\left(c_{\psi}\right)^{n+1}=0$.

These canonical cohomology classes of $E$ determine the characteristic number [19]

$$
\begin{equation*}
I_{\psi}=\int_{E} c_{1}(V T E) c_{\psi}^{n} \tag{1.2}
\end{equation*}
$$

$I_{\psi}$ depends only on the homotopy class of $\psi$. Moreover $I$ is an $\mathbb{R}$-valued group homomorphism on $\pi_{1}(\operatorname{Ham}(M, \omega))$, so the non-vanishing of $I$ implies that the group $\pi_{1}(\operatorname{Ham}(M, \omega))$ is infinite. That is, $I$ can be used to detect the infinitude of the corresponding homotopy group. Furthermore $I$ calibrates the Hofer's norm $v$ on $\pi_{1}(\operatorname{Ham}(M, \omega))$ in the sense that $v(\psi) \geq C\left|I_{\psi}\right|$, for all $\psi$, where $C$ is a positive constant [25].

In [28] we gave an explicit expression for the value of the characteristic number $I_{\psi}$. This value can be calculated if one has a family of local symplectic trivializations of $T M$ at one's disposal, whose domains cover $M$ and are fixed by the $\psi_{t}$ 's (see Theorem 3 in [28]). In this paper we use this theorem of [28] to prove that $\pi_{1}(\operatorname{Ham}(M))$ contains an infinite cyclic subgroup, when $M$ is the one point blow up of $\mathbb{C} P^{3}$. More precisely, in Section 3 we will prove the following result about the Hamiltonian group of the one point blow up of $\mathbb{C} P^{3}$.

Corollary 2. Let $(M, \omega)$ be the symplectic toric 6-manifold associated to the polytope obtained truncating the tetrahedron of $\mathbb{R}^{3}$ with vertices $(0,0,0),(\tau, 0,0),(0, \tau, 0),(0,0, \tau)$ by a horizontal plane [10], then $\pi_{1}(\operatorname{Ham}(M, \omega))$ contains an infinite cyclic subgroup.

Using quite different techniques McDuff and Tolman proved this result in [24].
We also give a sufficient condition for $\pi_{1}(\operatorname{Ham}(M))$ to contain a subgroup isomorphic to $\mathbb{Z}^{p}$, when $M$ is a general toric manifold. More precisely, let $\mathbb{T}$ be the torus $\left(S^{1}\right)^{r}$, and $\mathfrak{t}=\mathbb{R} \oplus \cdots \oplus \mathbb{R}$ its Lie algebra. Given $w_{j} \in \mathbb{Z}^{r}$, with $j=1, \ldots, m$ and $\tau \in \mathbb{R}^{r}$ we put

$$
\begin{equation*}
M=\left\{z \in \mathbb{C}^{m}: \pi \sum_{j=1}^{m}\left|z_{j}\right|^{2} w_{j}=\tau\right\} / \mathbb{T} \tag{1.3}
\end{equation*}
$$

where the relation defined by $\mathbb{T}$ is

$$
\begin{equation*}
\left(z_{j}\right) \simeq\left(z_{j}^{\prime}\right) \quad \text { iff } \quad \text { there is } \xi \in \mathfrak{t} \text { such that } z_{j}^{\prime}=z_{j} \mathrm{e}^{2 \pi \mathrm{i}\left(w_{j}, \xi\right\rangle} \text { for } j=1, \ldots, m \tag{1.4}
\end{equation*}
$$

We will assume that there is an open half space in $\mathbb{R}^{r}$ which contains all the vectors $w_{j}$ and that $\left\{w_{j}\right\}_{j}$ spans $\mathbb{R}^{r}$. We also assume that $\tau$ is a regular value of the map

$$
z \in \mathbb{C}^{m} \mapsto \pi \sum_{j=1}^{m}\left|z_{j}\right|^{2} w_{j} \in \mathbb{R}^{r}
$$

Then $M$ is a closed toric manifold of dimension $2 n:=2(m-r)$ [22].
For $j=1, \ldots, m$ we put ${ }^{j} \psi_{t}$ for the symplectomorphism

$$
{ }^{j} \psi_{t}:[z] \in M \mapsto\left[z_{1}, \ldots, z_{j} \mathrm{e}^{2 \pi \mathrm{i} t}, \ldots, z_{m}\right] \in M .
$$

Let $f_{j}$ be the corresponding normalized Hamiltonian and

$$
\begin{equation*}
\alpha_{j}:=\sum_{k=1}^{m}\left(\int_{\{[z]: z k=0\}} f_{j} \omega^{n-1}\right) . \tag{1.5}
\end{equation*}
$$

In Section 4 we will prove the following theorem.
Theorem 3. Let $(M, \omega)$ be the toric manifold defined by (1.3) and (1.4). If there are $p$ numbers linearly independent over $\mathbb{Z}$ in the set $\left\{\alpha_{1}, \ldots, \alpha_{m}\right\}$, where $\alpha_{j}$ is defined by (1.5), then $\pi_{1}(\operatorname{Ham}(M, \omega))$ contains a subgroup isomorphic to $\mathbb{Z}^{p}$. (If $p=0$ we mean by $\mathbb{Z}^{p}$ the trivial group.)

If $(M, \omega)$ is the toric manifold determined by the Delzant polytope $\Delta \subset \mathfrak{t}^{*}$, where $T$ is an $n$-dimensional torus, we deduce a formula for the value of $I$ on the Hamiltonian loops generated by the effective action of $T$ on $M$. In this formula geometrical magnitudes relative of $\Delta$ and the generator of the loop are involved (Proposition 9).

We denote by $M_{H}:=E \operatorname{Ham}(M) \times_{\operatorname{Ham}(M)} M \rightarrow B \operatorname{Ham}(M)$ the universal bundle, with fibre $M$, over the classifying space $B \operatorname{Ham}(M)$. Let $\mathbf{c} \in H^{2}\left(M_{H}, \mathbb{R}\right)$ denote the coupling class [14]. If $G$ is a compact Lie group and $\phi: G \rightarrow \operatorname{Ham}(M)$ be a group homomorphism, then $\phi$ induces a map $\Phi: B G \rightarrow B \operatorname{Ham}(M)$ between the corresponding classifying spaces. By means of this map the class $\mathbf{c}$ induces a class $c_{\phi} \in H^{2}\left(M_{\phi}\right)$, where $M_{\phi}=E G \times_{G} M$. The class $c_{\phi}$ is in fact the coupling class of the Hamiltonian fibration $M_{\phi} \xrightarrow{p} B G$ [21]. Using $c_{\phi}$ and $c_{1}^{\phi}(T M)$, the $G$-equivariant first Chern class of $T M$, we define

$$
R_{i}(\phi):=p_{*}\left(\left(c_{1}^{\phi}(T M)\right)^{i}\left(c_{\phi}\right)^{n-i}\right),
$$

$p_{*}$ being the integration along the fiber. The classes $R_{i}(\phi)$ were used in the paper [14] to study the cohomology of classifying spaces, and they are generalizations of the Miller-Morita-Mumford classes. Furthermore $R_{i}(\phi)$ can be calculated by the localization formula in $G$-equivariant cohomology.

In the set of all Lie group homomorphisms from $G$ to $\operatorname{Ham}(M)$ we declare that two elements are equivalent if they are homotopic by means of a family of Lie group homomorphisms from $G$ to $\operatorname{Ham}(M)$. The quotient set is denoted by $[G, \operatorname{Ham}(M)]_{g h}$. If $\phi$ and $\tilde{\phi}$ are two Hamiltonian $G$-actions on $M$ which define the same element in $[G, \operatorname{Ham}(M)]_{g h}$, then $R_{i}(\phi)=R_{i}(\tilde{\phi})$. Thus, one has a numerical criterion for two Hamiltonian $G$-actions not to be equivalent under homotopies consisting of Hamiltonian $G$-actions. We will prove the existence of pairs of Hamiltonian circle actions in a Hirzebruch surface, which define the same element in $\pi_{1}(\mathrm{Ham})$ but its classes in $[U(1), \operatorname{Ham}(M)]_{g h}$ are not equal.

If $\mu: M \rightarrow \mathfrak{g}^{*}$ is a moment map for a $G$-action $\phi$ on $M$, we denote with $\Omega_{\phi}$ the equivariant closed 2-form $\omega+\mu$ [13]. Given $Z \in \mathfrak{g}$ we denote by $\mathcal{R}_{\phi}(Z)$ the number

$$
\begin{equation*}
\mathcal{R}_{\phi}(Z):=(2 \pi \mathrm{i})^{-n} \int_{M} \mathrm{e}^{\mathrm{i} \Omega_{\phi}(Z)} . \tag{1.6}
\end{equation*}
$$

When $G=U(1)$ and $\mu$ is normalized, $\Omega_{\phi}$ is a representative of the coupling class $c_{\phi}$. If $X \in \mathfrak{g}$ generates a loop $\phi$ in $\operatorname{Ham}(\mathcal{O})$, where $\mathcal{O}$ is a regular coadjoint orbit of $G$, then $\mathcal{R}_{\phi}(X)$ is the Fourier transform of the orbit and the Harish-Chandra Theorem (see [5]) allows us to calculate $\mathcal{R}_{\phi}(X)$ in terms of the root structure of $\mathfrak{g}$. Using this fact we give an example of two Hamiltonian circle actions on $\mathbb{C} P^{1}$ which define the same element in $\pi_{1}(\mathrm{Ham})$ but they are not homotopic by a family of circle actions.

The paper is organized as follows. In Section 2 we calculate the Maslov index of the linearized flow of certain loops in the Hamiltonian group of flag manifolds, and determine lower bound for the corresponding $\sharp \pi_{1}(\mathrm{Ham})$.

If $M$ is the one point blow up of $\mathbb{C} P^{3}$, then $M$ is a submanifold of $\mathbb{C} P^{1} \times \mathbb{C} P^{3}$, and the natural actions of $U(1)$ on the projective spaces define Hamiltonian loops $\psi$ in $M$. Section 3 is concerned with the determination of $I_{\psi}$ for these loops. As a consequence of these calculations we deduce that $\mathbb{Z} \subset \pi_{1}(\operatorname{Ham}(M))$.

In Section 4 we generalize the arguments developed in Section 3 to toric manifolds. From this generalization it follows a sufficient condition, stated in Theorem 3, for the existence of an infinite subgroup in $\pi_{1}(\operatorname{Ham}(M))$, when $M$ is a toric manifold. Finally we check that this sufficient condition does not hold for $\mathbb{C} P^{n}$ with $n=1,2$. This is consistent with the fact that $\pi_{1}\left(\operatorname{Ham}\left(\mathbb{C} P^{n}\right)\right)$ is finite for $n=1,2$.

In Section 5 the $R_{i}(\phi)$ are introduced. If $M$ is a Hirzebruch surface, the toric structure allows us to define two Hamiltonian $S^{1}$-actions $\phi, \tilde{\phi}$, which are not homotopic by means of a set of $U(1)$-actions, since $R_{1}(\phi) \neq R_{1}(\tilde{\phi})$. When $M$ satisfies some additional hypotheses we will prove the existence of infinitely many pairs ( $\zeta, \zeta^{\prime}$ ) of circle actions on $M$, such that $[\zeta]=\left[\zeta^{\prime}\right] \in \pi_{1}(\operatorname{Ham}(M))$, but $[\zeta] \neq\left[\zeta^{\prime}\right] \in[U(1) \text {, } \operatorname{Ham}(M)]_{g h}$.

Conventions. We use the following conventions. If $f_{t}$ is a time-dependent Hamiltonian on $(M, \omega)$, the corresponding Hamiltonian vector $Y_{t}$ is defined by

$$
\begin{equation*}
\iota_{Y_{t}} \omega=-\mathrm{d} f_{t} . \tag{1.7}
\end{equation*}
$$

This time-dependent vector field vector determines the respective family $\psi_{t}$ of symplectomorphisms by

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} \psi_{t}=Y_{t} \circ \psi_{t}, \quad \psi_{0}=\mathrm{id} \tag{1.8}
\end{equation*}
$$

If the group $G$ acts on $M$ and $X \in \mathfrak{g}$, by $X_{M}$ is denoted the vector field whose value at $x \in M$ is

$$
\begin{equation*}
X_{M}(x)=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \mathrm{e}^{-t X} x \tag{1.9}
\end{equation*}
$$

If the action of $G$ is Hamiltonian, a moment map $\mu: M \rightarrow \mathfrak{g}^{*}$ satisfies $\mathrm{d} \mu(X)=\iota_{X_{M}} \omega$. According to (1.7) and (1.8) the function $f:=\mu(X)$ defines the isotopy $\phi_{t}^{X}$ given by

$$
\begin{equation*}
\phi_{t}^{X}(x):=\mathrm{e}^{t X} x \tag{1.10}
\end{equation*}
$$

Given $E \rightarrow M$ is a $G$-equivariant bundle, $s$ is a section of $E$ and $X \in \mathfrak{g}$, we define the action of $X$ on $s$ as the section $X s$

$$
\begin{equation*}
(X s)(x)=\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0}\left(\mathrm{e}^{t X}\right) \cdot s\left(\mathrm{e}^{-t X} x\right) \tag{1.11}
\end{equation*}
$$

## 2. Lower bounds for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O}))$

### 2.1. Maslov index of the linearized flow

We denote by $(M, \omega)$ a closed, connected, symplectic, $2 n$-dimensional manifold. Let $\psi: \mathbb{R} / \mathbb{Z} \rightarrow \operatorname{Ham}(M, \omega)$ be a loop in the group of Hamiltonian symplectomorphisms at Id. Given $x \in M$, the curve $C:=\left\{\psi_{t}(x) \mid t \in[0,1]\right\}$ is null-homotopic [18]. Let $S$ be a 2-dimensional singular disc in $M$ whose boundary is $C$, and let $X_{1}, \ldots, X_{2 n}$ be vector fields on $S$ which form a symplectic basis of $T_{p} M$ for each $p \in S$. Then

$$
\left(\psi_{t}\right)_{*}\left(X_{i}(x)\right)=\sum_{k} A_{i}^{k}(t, x) X_{k}\left(x_{t}\right),
$$

with $x_{t}:=\psi_{t}(x)$ and $A \in \operatorname{Sp}(2 n, \mathbb{R})$. By $\rho$ will be denoted the usual map $\operatorname{Sp}(2 n, \mathbb{R}) \rightarrow U(1)$ which restricts to the determinant map on $U(n)$ [26]. Setting $a(t, x):=\rho(A(t, x))$, we write $J_{\psi}(X, x)$ for the winding number of the map $t \in \mathbb{R} / \mathbb{Z} \rightarrow a(t, x) \in U(1)$. That is,

$$
J_{\psi}(X, x)=\frac{1}{2 \pi \mathrm{i}} \int_{0}^{1} a^{-1} \frac{\partial a}{\partial t}(t, x) \mathrm{d} t .
$$

If $N$ is the minimal Chern number of $M$ on spheres, the class of $J_{\psi}(X, x)$ in $\mathbb{Z} / 2 N \mathbb{Z}$ only depends on the homotopy class of $[\psi]$. The element in $\mathbb{Z} / 2 N \mathbb{Z}$ defined by $J_{\psi}(X, x)$ will be denoted $J[\psi]$ and is the Maslov index of the flow $\psi_{t *}$.

### 2.2. Coadjoint orbits

Let $G$ be a compact semisimple Lie group, and $\eta$ an element of $\mathfrak{g}^{*}$, the dual of the Lie algebra of $G$. We denote by $\mathcal{O}$ the coadjoint orbit of $\eta$ equipped with the standard symplectic structure [17]. This orbit can be identified with $G / G_{\eta}$, where $G_{\eta}$ is the stabilizer of $\eta$ for the coadjoint action of $G$. The subgroup $G_{\eta}$ contains a maximal torus $T$ of $G$ [12]. We have the decomposition of $\mathfrak{g}_{\mathbb{C}}$ as a direct sum of root spaces

$$
\mathfrak{g}_{\mathbb{C}}=\mathfrak{t}_{\mathbb{C}} \oplus \bigoplus_{\alpha \in \Phi} \mathfrak{g}_{\alpha}
$$

with $\Phi$ the set of roots determined by $T$. We denote by $\check{\alpha} \in\left[\mathfrak{g}_{\alpha}, \mathfrak{g}_{-\alpha}\right]$ the coroot of $\alpha$. Let $\mathfrak{p}$ be the parabolic subalgebra

$$
\mathfrak{p}=\mathfrak{t}_{\mathbb{C}} \oplus \bigoplus_{\eta(i \tilde{\alpha}) \geq 0} \mathfrak{g}_{\alpha}
$$

By $Z_{A}$ is denoted the right invariant vector field on $G$ determined by $A \in \mathfrak{p}$. Since $\mathfrak{p}$ is a subalgebra, $\left\{Z_{A} \mid A \in \mathfrak{p}\right\}$ defines an integrable distribution on $G$. Its projection onto $G / G_{\eta}$ is a complex structure on the orbit $\mathcal{O}$ compatible with the symplectic structure. If $P$ is the parabolic subgroup of $G_{\mathbb{C}}$ generated by $\mathfrak{p}, G_{\mathbb{C}} / P$ is this complexification of $G / G_{\eta}=\mathcal{O}$, and

$$
T_{\eta}^{1,0} \simeq \mathfrak{g}_{\mathbb{C}} / \mathfrak{p} \simeq \bigoplus_{\alpha \in \Lambda} \mathfrak{g}_{\alpha}=: \mathfrak{n}
$$

where $\Lambda=\left\{\beta_{1}, \ldots, \beta_{r}\right\}$ is a subset of $\Phi$. Let $A_{1}, \ldots, A_{r}$ be a $\mathbb{C}$-basis for $\mathfrak{n}$, with $A_{j} \in \mathfrak{g}_{\beta_{j}}$, then $\left\{Z_{A_{j}}\right\}_{j}$ is a local frame for $T^{1,0} \mathcal{O}$ on a neighborhood $U$ of $\eta$.

If $\left\{g_{t} \mid t \in[0,1]\right\}$ is a family of elements of $G$, such that $g_{0}=e$ and $g_{1} \in Z(G)$, then

$$
\left\{\psi_{t}: g G_{\eta} \in G / G_{\eta} \mapsto g_{t} g G_{\eta} \in G / G_{\eta}\right\}_{t \in[0,1]}
$$

is a loop in $\operatorname{Ham}(\mathcal{O})$. Furthermore

$$
\begin{equation*}
\left(\psi_{t}\right)_{*} Z_{A_{j}}=Z_{g_{t} \cdot A_{j}} \tag{2.1}
\end{equation*}
$$

with $g \cdot A:=\operatorname{Ad}_{g} A$.
Let $g_{t}=\exp \left(C_{t}\right)$, with $C_{t} \in \mathfrak{t}$ and $C_{0}=0$. As $\left[C_{t}, E\right]=\alpha\left(C_{t}\right) E$ if $E \in \mathfrak{g}_{\alpha}$, then we have

$$
\operatorname{Ad}_{g_{t}} A_{j}=\exp \left(\beta_{j}\left(C_{t}\right)\right) A_{j}
$$

It follows from (2.1) that for $v \in U$ the matrix of $\left(\psi_{t}\right)_{*}(\nu)$ with respect $\left\{Z_{A_{j}}\right\}_{j}$ is

$$
\operatorname{diag}\left(\mathrm{e}^{\beta_{1}\left(C_{t}\right)}, \ldots, \mathrm{e}^{\beta_{r}\left(C_{t}\right)}\right) \in U(r)
$$

The Maslov index $J[\psi]$ is the class in $\mathbb{Z} / 2 N \mathbb{Z}$ of the winding number of the map

$$
\begin{equation*}
t \in[0,1] \mapsto \exp \left(\sum_{j=1}^{r} \beta_{j}\left(C_{t}\right)\right)=\exp \left(\sum_{\alpha \in \Lambda} \alpha\left(C_{t}\right)\right) \in U(1) . \tag{2.2}
\end{equation*}
$$

That is,

$$
J[\psi]=\frac{1}{2 \pi \mathrm{i}} \sum_{\alpha \in \Lambda} \alpha\left(C_{1}\right)+2 N \mathbb{Z} .
$$

We have the following proposition.

Proposition 4. Let $\mathcal{O}$ be a coadjoint orbit of $G$ whose complexification is $G_{\mathbb{C}} / P$, with $\mathfrak{g}_{\mathbb{C}} / \mathfrak{p} \simeq \bigoplus_{\alpha \in \Lambda} \mathfrak{g}_{\alpha}$. Let $\left\{C_{t}\right\}_{t \in[0,1]}$ be a curve in $\mathfrak{t}$ such that $C_{0}=0$ and $\exp \left(C_{1}\right) \in Z(G)$. If

$$
\frac{1}{2 \pi \mathrm{i}} \sum_{\alpha \in \Lambda} \alpha\left(C_{1}\right) \notin 2 N \mathbb{Z}
$$

then $g_{t}=\exp \left(C_{t}\right)$ defines a nontrivial element in $\pi_{1}(\operatorname{Ham}(\mathcal{O}))$.

### 2.3. Flag manifolds in $\mathbb{C}^{n+1}$

From now to the end of Section $2 G$ will be the group $S U(n+1)$, and $T$ the subgroup of diagonal elements. We denote by $\Delta$ the usual base of roots; that is, $\Delta=\left\{\alpha_{1}, \ldots, \alpha_{n}\right\}$, where $\alpha_{i}=\epsilon_{i}-\epsilon_{i+1}$ (we use the notation of [8]). Each subset $I \subset \Delta$ determines a parabolic subgroup $P_{I}$ of $S L(n+1, \mathbb{C})$. This subgroup is generated by the subalgebra

$$
\mathfrak{p}_{I}=\mathfrak{t}_{\mathbb{C}} \oplus \bigoplus_{\alpha \in \tilde{I}} \mathfrak{g}_{\alpha}
$$

where $\tilde{I}$ consists of all roots that can be written as sums of negative elements in $I$ together with all positive roots [7]. If $I=\Delta-\left\{\alpha_{n}\right\}$, then $\mathfrak{p}_{I}=\mathfrak{g l} l(n, \mathbb{C})$ and $\mathfrak{s l}(n+1, \mathbb{C}) / \mathfrak{p}_{I}$ is isomorphic

$$
\bigoplus_{j=1}^{n} \mathfrak{g}_{\beta_{j}}=\mathfrak{n}
$$

with $\beta_{j}=\epsilon_{n+1}-\epsilon_{j}$. In this case

$$
S L(n+1, \mathbb{C}) / P_{I} \simeq S U(n+1) / U(n)=\mathbb{C} P^{n} .
$$

Next we will determine a lower bound for $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O}))$, when $\mathcal{O}$ is a coadjoint orbit of $S U(n+1)$ diffeomorphic to $\mathbb{C} P^{n}$. Let us take a complex number $z$ such that $z^{n+1}=1$, and put

$$
g_{t}:=\operatorname{diag}\left(z^{t}, \ldots, z^{t}, z^{-n t}\right) \in T \subset S U(n+1) .
$$

So $g_{1} \in Z(S U(n+1))$, and moreover $g_{t}=\exp \left(C_{t}\right)$, with

$$
\begin{equation*}
C_{t}=\frac{2 k \pi \mathrm{i} t}{n+1} \operatorname{diag}(1, \ldots, 1,-n) \tag{2.3}
\end{equation*}
$$

where $k$ is any element of $\{0,1, \ldots, n\}$. Then $\left(\epsilon_{n+1}-\epsilon_{j}\right)\left(C_{t}\right)=-2 k \pi t \mathrm{i}$, and in this case the map (2.2) is

$$
t \in[0,1] \mapsto \exp (-2 k n \pi t \mathrm{i}) \in U(1)
$$

whose winding number is $-k n$. Hence, for $k=0,1, \ldots, n$ we obtain loops $\left\{_{k} \psi_{t} \mid t \in[0,1]\right\}$ in $\operatorname{Ham}(\mathcal{O})$ such that the corresponding Maslov indices take the values

$$
J\left[{ }_{k} \psi\right]=-k n+2 N \mathbb{Z}
$$

The minimal Chern number of $\mathbb{C} P^{n}$ is equal to $n+1$. As $-k n+2(n+1) \mathbb{Z} \neq-j n+2(n+1) \mathbb{Z}$ for $k \neq j \in\{0,1, \ldots n\}$, then $\left[{ }_{k} \psi\right] \neq\left[{ }_{j} \psi\right] \in \pi_{1}(\operatorname{Ham}(\mathcal{O}))$. We have proved the following theorem.

Theorem 5. If $\mathcal{O}$ is a coadjoint orbit of $S U(n+1)$ diffeomorphic to $\mathbb{C} P^{n}$, then $\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O})) \geq n+1$.
It is known that $\pi_{1}\left(\operatorname{Ham}\left(\mathbb{C} P^{1}\right)\right)=\mathbb{Z} / 2 \mathbb{Z}$ and that $\operatorname{Ham}\left(\mathbb{C} P^{2}\right)$ has the homotopy type of $P U(3)$ [9], so $\sharp \pi_{1}\left(\operatorname{Ham}\left(\mathbb{C} P^{2}\right)\right)=3$. The bound given in Theorem 5 is compatible with those facts.
Proof of Theorem 1. Now we consider coadjoint orbits of $S U(n+1)$ which are diffeomorphic to the Grassmannian $G_{s}\left(C^{n+1}\right)$ of $s$-dimensional subspaces of $\mathbb{C}^{n+1}$. Let $\mathfrak{p}$ be the parabolic subalgebra generated by $I=\Delta-\left\{\alpha_{s}\right\}$; that is, we delete the $s$-node in the Dynkin diagram. (If $s=n$, the corresponding Grassmannian is $\mathbb{C} P^{n}$.) Now

$$
\mathfrak{s l}(n+1, \mathbb{C}) / \mathfrak{p}=\bigoplus_{\beta} \mathfrak{g}_{\beta}
$$

with $\beta=\epsilon_{j}-\epsilon_{i}, j=s+1, \ldots, n+1$ and $i=1, \ldots, s$.

With the above notations $\left(\epsilon_{j}-\epsilon_{i}\right)\left(C_{t}\right)=0$ for any $i, j$ with $j \neq n+1$. Now the map (2.2) is

$$
t \in[0,1] \mapsto \exp (-2 k s t \pi \mathrm{i}) \in U(1)
$$

and its winding number is $-k s$. The minimal Chern number $N$ for the Grassmannian $G_{s}\left(\mathbb{C}^{n+1}\right)$ is $n+1$. If $s$ and $n+1$ are relatively prime then

$$
\sharp\{-k s+2(n+1) \mathbb{Z} \mid k=0,1, \ldots n\}=n+1 .
$$

Given $\mathfrak{p}$ a parabolic subalgebra of $\mathfrak{l l}(n+1, \mathbb{C})$ which contains the standard Borel subalgebra, then

$$
\mathfrak{s l}(n+1, \mathbb{C}) / \mathfrak{p} \simeq \bigoplus_{\beta \in \Lambda} \mathfrak{g}_{\beta}
$$

where $\Lambda=\Phi \backslash \tilde{I}$.
Given $a \in\{1, \ldots, n+1\}$ we put

$$
\begin{equation*}
\langle a\rangle=\sharp\left\{\beta=\epsilon_{i}-\epsilon_{a} \in \Lambda\right\}-\sharp\left\{\beta=\epsilon_{a}-\epsilon_{j} \in \Lambda\right\} . \tag{2.4}
\end{equation*}
$$

Let $C_{t}(a)$ be the element of $\mathfrak{t}$ defined by

$$
\begin{equation*}
C_{t}(a)=\frac{2 k \pi \mathrm{i} t}{n+1} \operatorname{diag}(1, \ldots, 1,-n, 1, \ldots, 1) \tag{2.5}
\end{equation*}
$$

where $-n$ is in the position $a$. The element $C_{t}$ in (2.3) is equal to $C_{t}(n+1)$. We consider the curve $g_{t}=\exp \left(C_{t}(a)\right)$, then

$$
\sum_{\beta \in \Lambda} \beta\left(C_{t}(a)\right)=2 k \mathrm{i} t\langle a\rangle,
$$

and the winding number of the map $t \mapsto \exp \sum \beta\left(C_{t}(a)\right)$ is $k\langle a\rangle$. Hence

$$
\sharp \pi_{1}(\operatorname{Ham}(S L(n+1, \mathbb{C}) / P)) \geq \sharp\{k\langle a\rangle+2 N \mathbb{Z} \mid k=0,1, \ldots, n\} .
$$

So one arrives at the following result.
Theorem 6. If $\mathcal{O}$ is a coadjoint orbit of $S U(n+1)$ diffeomorphic to the flag manifold $S L(n+1, \mathbb{C}) / P$, then

$$
\sharp \pi_{1}(\operatorname{Ham}(\mathcal{O})) \geq \max _{a=1, \ldots, n+1}(\sharp\{k\langle a\rangle+2 N \mathbb{Z} \mid k=0,1, \ldots, n\}),
$$

where the integer $\langle a\rangle$ is defined by the parabolic subalgebra $\mathfrak{p}$ by (2.4).

## 3. Hamiltonian group of the one point blow up of $\mathbb{C} P^{3}$

Given $\tau, \sigma \in \mathbb{R}_{>0}$, with $\sigma<\tau$, let $M$ be the following manifold

$$
\begin{equation*}
M=\left\{z \in \mathbb{C}^{5}:\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}+\left|z_{5}\right|^{2}=\tau / \pi,\left|z_{3}\right|^{2}+\left|z_{4}\right|^{2}=\sigma / \pi\right\} / \mathbb{T}, \tag{3.1}
\end{equation*}
$$

where the action of $\mathbb{T}=\left(S^{1}\right)^{2}$ is defined by

$$
\begin{equation*}
(a, b)\left(z_{1}, z_{2}, z_{3}, z_{4}, z_{5}\right)=\left(a z_{1}, a z_{2}, a b z_{3}, b z_{4}, a z_{5}\right) \tag{3.2}
\end{equation*}
$$

for $a, b \in S^{1}$.
$M$ is a toric 6-manifold; more precisely, it is the toric manifold associated to the polytope obtained truncating the tetrahedron of $\mathbb{R}^{3}$ with vertices

$$
(0,0,0),(\tau, 0,0),(0, \tau, 0),(0,0, \tau)
$$

by a horizontal plane through the point $(0,0, \lambda)$, with $\lambda:=\tau-\sigma[10]$.

For $0 \neq z_{j} \in \mathbb{C}$ we put $z_{j}=\rho_{j} \mathrm{e}^{\mathrm{i} \theta_{j}}$, with $\left|z_{j}\right|=\rho_{j}$. On the set of points $[z] \in M$ with $z_{i} \neq 0$ for all $i$ one can consider the coordinates

$$
\begin{equation*}
\left(\frac{\rho_{1}^{2}}{2}, \varphi_{1}, \frac{\rho_{2}^{2}}{2}, \varphi_{2}, \frac{\rho_{3}^{2}}{2}, \varphi_{3}\right) \tag{3.3}
\end{equation*}
$$

where the angle coordinates are defined by

$$
\begin{equation*}
\varphi_{1}=\theta_{1}-\theta_{5}, \quad \varphi_{2}=\theta_{2}-\theta_{5}, \quad \varphi_{3}=\theta_{3}-\theta_{4}-\theta_{5} . \tag{3.4}
\end{equation*}
$$

Then the standard symplectic structure on $\mathbb{C}^{5}$ induces the following form $\omega$ on this part of $M$

$$
\begin{equation*}
\omega=\sum_{j=1}^{3} \mathrm{~d}\left(\frac{\rho_{j}^{2}}{2}\right) \wedge \mathrm{d} \varphi_{j} . \tag{3.5}
\end{equation*}
$$

### 3.1. Darboux coordinates on $M$

Let $0<\epsilon \ll 1$, we write

$$
B_{0}=\left\{[z] \in M:\left|z_{j}\right|>\epsilon, \text { for all } j\right\} .
$$

For each $j \in\{1,2,3,4,5\}$ we set

$$
B_{j}=\left\{[z] \in M:\left|z_{j}\right|<2 \epsilon \text { and }\left|z_{i}\right|>\epsilon, \text { for all } i \neq j\right\} .
$$

The family $B_{0}, \ldots, B_{5}$ is not a covering of $M$, but if $[z] \notin \cup B_{k}$, then there are $i, j$, with $i \neq j$ and $\left|z_{i}\right| \leq \epsilon \geq\left|z_{j}\right|$.
We will define Darboux coordinates on $B_{0}, \ldots, B_{5}$. On $B_{0}$ we will consider the well-defined Darboux coordinates (3.3).

On $B_{1}, \rho_{j} \neq 0$ for $j \neq 1$; so the angle coordinates $\varphi_{2}$ and $\varphi_{3}$ of (3.4) are well-defined. We define $x_{1}, y_{1}$ by the relation $x_{1}+\mathrm{i} y_{1}:=\rho_{1}{ }^{\mathrm{i} \varphi_{1}}$ and $x_{1}=0=y_{1}$, if $z_{1}=0$. In this way we take as symplectic coordinates on $B_{1}$

$$
\left(x_{1}, y_{1}, \frac{\rho_{2}^{2}}{2}, \varphi_{2}, \frac{\rho_{3}^{2}}{2}, \varphi_{3}\right) .
$$

We will also consider the following Darboux coordinates: On $B_{2}$

$$
\left(\frac{\rho_{1}^{2}}{2}, \varphi_{1}, x_{2}, y_{2}, \frac{\rho_{3}^{2}}{2}, \varphi_{3}\right), \quad \text { with } x_{2}+\mathrm{i} y_{2}:=\rho_{2} \mathrm{e}^{\mathrm{i} \varphi_{2}} ; \text { and } x_{2}=0=y_{2}, \text { if } z_{2}=0
$$

On $B_{3}$

$$
\left(\frac{\rho_{1}^{2}}{2}, \varphi_{1}, \frac{\rho_{2}^{2}}{2}, \varphi_{2}, x_{3}, y_{3}\right), \quad \text { where } x_{3}+\mathrm{i} y_{3}:=\rho_{3} \mathrm{e}^{\mathrm{i} \varphi_{3}} .
$$

On $B_{4}$

$$
\left(\frac{\rho_{1}^{2}}{2}, \varphi_{1}, \frac{\rho_{2}^{2}}{2}, \varphi_{2}, x_{4}, y_{4}\right), \quad \text { with } x_{4}+\mathrm{i} y_{4}:=\rho_{4} \mathrm{e}^{\mathrm{i} \varphi_{4}} \text { and } \varphi_{4}=\theta_{4}-\theta_{3}+\theta_{5} .
$$

On $B_{5}$

$$
\left(x_{5}, y_{5}, \frac{\rho_{2}^{2}}{2}, \chi_{2}, \frac{\rho_{3}^{2}}{2}, \chi_{3}\right),
$$

where

$$
x_{5}+\mathrm{i} y_{5}:=\rho_{5} \mathrm{e}^{\mathrm{i} \chi_{5}}, \quad \chi_{2}=\theta_{2}-\theta_{1}, \quad \chi_{3}=\theta_{3}-\theta_{1}-\theta_{4}, \quad \chi_{5}=\theta_{5}-\theta_{1} .
$$

If $\left[z_{1}, \ldots, z_{5}\right]$ is a point of

$$
M \backslash \bigcup_{i=0}^{5} B_{i}
$$

then there are $a \neq b \in\{1, \ldots, 5\}$ such that $\left|z_{a}\right|,\left|z_{b}\right| \leq \epsilon$. We can cover the set $M \backslash \bigcup B_{i}$ by Darboux charts denoted $B_{6}, \ldots, B_{q}$ similar to the preceding $B_{i}$ 's satisfying the following condition. The image of each $B_{a}$, with $a=6, \ldots, q$, is contained in a prism of $\mathbb{R}^{6}$ of the form

$$
\prod_{i=1}^{6}\left[c_{i}, d_{i}\right]
$$

where at least two intervals $\left[c_{i}, d_{i}\right]$ have length of order $\epsilon$.
By the infinitesimal "size" of the $B_{j}$, for $j \geq 1$, it turns out

$$
\begin{equation*}
\int_{B_{j}} \omega^{3}=O(\epsilon), \quad \text { for } j \geq 1 \tag{3.6}
\end{equation*}
$$

### 3.2. A loop in $\operatorname{Ham}(M)$

Let $\psi_{t}$ be the symplectomorphism of $M$ defined by

$$
\begin{equation*}
\psi_{t}[z]=\left[z_{1} \mathrm{e}^{2 \pi \mathrm{i} t}, z_{2}, z_{3}, z_{4}, z_{5}\right] \tag{3.7}
\end{equation*}
$$

Then $\left\{\psi_{t}\right\}_{t}$ is a loop in the group $\operatorname{Ham}(M)$ of Hamiltonian symplectomorphisms of $M$. By $f$ is denoted the corresponding normalized Hamiltonian function. Hence $f=\pi \rho_{1}^{2}-\kappa$ with $\kappa \in \mathbb{R}$ such that $\int_{M} f \omega^{3}=0$.

We will calculate $I_{\psi}$ using the following result proved in [28] (Theorem 3 of [28]).
Theorem 7. Let $\psi: S^{1} \rightarrow \operatorname{Ham}(M, \omega)$ be a closed Hamiltonian isotopy generated by the normalized timedependent Hamiltonian $f_{t}$. If $\left\{B_{1}, \ldots, B_{m}\right\}$ is a set of symplectic trivializations for $T M$ which covers $M$ and such that $\psi_{t}\left(B_{j}\right)=B_{j}$, for all $t$ and all $j$, then

$$
\begin{equation*}
I_{\psi}=\sum_{i=1}^{m} J_{i} \int_{B_{i} \backslash \bigcup_{j<i} B_{j}} \omega^{n}+\sum_{i<k} N_{i k}, \tag{3.8}
\end{equation*}
$$

where

$$
N_{i k}=n \frac{\mathrm{i}}{2 \pi} \int_{S^{1}} \mathrm{~d} t \int_{A_{i k}}\left(f_{t} \circ \psi_{t}\right)\left(\mathrm{d} \log r_{i k}\right) \wedge \omega^{n-1}
$$

$A_{i k}=\left(\partial B_{i} \backslash \cup_{r<k} B_{r}\right) \cap B_{k}, J_{i}$ is the Maslov index of $\left(\psi_{t}\right)_{*}$ in the trivialization $B_{i}$ and $r_{i k}$ the corresponding transition function of $\operatorname{det}(T M)$.

We will prove that, in the case we are considering, some summands in (3.8) are of order $\epsilon$. We will neglect the order $\epsilon$ summands, and in this way we will obtain an expression which is equal to $I_{\psi}$ up to an addend of order $\epsilon$.

In the coordinates (3.3) of $B_{0}, \psi_{t}$ is the map $\varphi_{1} \mapsto \varphi_{1}+2 \pi t$. So the Maslov index $J_{B_{0}}=0$. It follows from (3.6) and Theorem 7

$$
\begin{equation*}
I_{\psi}=\sum_{i<k} N_{i k}+O(\epsilon), \tag{3.9}
\end{equation*}
$$

with

$$
N_{i k}=\frac{3 \mathrm{i}}{2 \pi} \int_{A_{i k}} f \mathrm{~d} \log r_{i k} \wedge \omega^{2}
$$

If $[z] \in A_{i k} \subset \partial B_{i} \cap B_{k}$, with $1 \leq i<k$, then at least the modules $\left|z_{a}\right|$ and $\left|z_{b}\right|$ of two components of $[z]$ are of order $\epsilon$; so $N_{i k}$ is of order $\epsilon$ when $1 \leq i<k$. Analogously $N_{0 k}$ is of order $\epsilon$, for $k=6, \ldots, q$. Hence (3.9) reduces to

$$
\begin{equation*}
I_{\psi}=\sum_{k=1}^{5} N_{0 k}+O(\epsilon) \tag{3.10}
\end{equation*}
$$

If we put

$$
\begin{equation*}
N_{0 k}^{\prime}=\frac{3 \mathrm{i}}{2 \pi} \int_{A_{0 k}^{\prime}} f \mathrm{~d} \log r_{i k} \wedge \omega^{2} \tag{3.11}
\end{equation*}
$$

with

$$
A_{0 k}^{\prime}=\left\{[z] \in M:\left|z_{k}\right|=\epsilon,\left|z_{r}\right|>\epsilon \text { for all } r \neq k\right\}
$$

then

$$
N_{0 k}=N_{0 k}^{\prime}+O(\epsilon)
$$

and

$$
\begin{equation*}
I_{\psi}=\sum_{k=1}^{5} N_{0 k}^{\prime}+O(\epsilon) \tag{3.12}
\end{equation*}
$$

### 3.3. Calculation of the $N_{0 k}^{\prime}$ 's

First we determine the value of $N_{01}^{\prime}$. To know the transition function $r_{01}$ one needs the Jacobian matrix $R$ of the transformation

$$
\left(x_{1}, y_{1}, \frac{\rho_{2}^{2}}{2}, \varphi_{2}, \frac{\rho_{3}^{2}}{2}, \varphi_{3}\right) \rightarrow\left(\frac{\rho_{1}^{2}}{2}, \varphi_{1}, \frac{\rho_{2}^{2}}{2}, \varphi_{2}, \frac{\rho_{3}^{2}}{2}, \varphi_{3}\right)
$$

in the points of $A_{01}^{\prime}$; where $\rho_{1}^{2}=x_{1}^{2}+y_{1}^{2}, \varphi_{1}=\tan ^{-1}\left(y_{1} / x_{1}\right)$. The function $r_{01}=\rho(R)$, where $\rho: \operatorname{Sp}(6, \mathbb{R}) \rightarrow U(1)$ is the map which restricts to the determinant on $U(3)$ [26]. The non-trivial block of $R$ is the diagonal one

$$
\left(\begin{array}{cc}
x_{1} & y_{1} \\
r & s
\end{array}\right)
$$

with $r=-y_{1}\left(x_{1}^{2}+y_{1}^{2}\right)^{-1}$ and $s=x_{1}\left(x_{1}^{2}+y_{1}^{2}\right)^{-1}$. The non-real eigenvalues of $R$ are

$$
\lambda_{ \pm}=\frac{x_{1}+s}{2} \pm \frac{\mathrm{i} \sqrt{4-\left(s+x_{1}\right)^{2}}}{2}
$$

These non-real eigenvalues occur when $\left(s+x_{1}\right)^{2}<2$. On $A_{01}^{\prime}$ this condition is equivalent to $\left|\cos \varphi_{1}\right|<2 \epsilon\left(\epsilon^{2}+\right.$ 1) ${ }^{-1}=: \delta$, since $\rho_{1}=\epsilon$ for the points of $A_{01}^{\prime}$.

If $y_{1}>0$ then $\lambda_{-}$of the first kind (see [26]) and $\lambda_{+}$is of the first kind if $y_{1}<0$. Hence, on $A_{01}^{\prime}$,

$$
\rho(R)= \begin{cases}\lambda_{+}\left|\lambda_{+}\right|^{-1}=x+\mathrm{i} y, & \text { if }\left|\cos \varphi_{1}\right|<\delta \text { and } y_{1}<0 \\ \lambda_{-}\left|\lambda_{-}\right|^{-1}=x-\mathrm{i} y, & \text { if }\left|\cos \varphi_{1}\right|<\delta \text { and } y_{1}>0 \\ \pm 1, & \text { otherwise }\end{cases}
$$

where $x=\delta^{-1} \cos \varphi_{1}$, and $y=\sqrt{1-x^{2}}$.
If we put $\rho(R)=\mathrm{e}^{\mathrm{i} \gamma}$ then, for the points of $A_{01}^{\prime}$ in which $\left|\cos \varphi_{1}\right|<\delta$,

$$
\cos \gamma=\delta^{-1} \cos \varphi_{1}, \quad \text { and } \quad \sin \gamma= \begin{cases}-\sqrt{1-\cos ^{2} \gamma}, & \text { if } \sin \varphi_{1}>0 \\ \sqrt{1-\cos ^{2} \gamma}, & \text { if } \sin \varphi_{1}<0\end{cases}
$$

So, when $\varphi_{1}$ runs anticlockwise from 0 to $2 \pi, \gamma$ goes round the circumference clockwise; that is, $\gamma=h\left(\varphi_{1}\right)$, where $h$ is a function such that

$$
\begin{equation*}
h(0)=2 \pi, \quad \text { and } \quad h(2 \pi)=0 . \tag{3.13}
\end{equation*}
$$

As $r_{01}=\rho(R)$, then $\mathrm{d} \log r_{01}=\mathrm{id} h$.
On $A_{01}^{\prime}$ the symplectic form (3.5) reduces to (1/2) $\left(\mathrm{d} \rho_{2}^{2} \wedge \mathrm{~d} \varphi_{2}+\mathrm{d} \rho_{3}^{2} \wedge \mathrm{~d} \varphi_{3}\right)$. From (3.11) one deduces

$$
\begin{equation*}
N_{01}^{\prime}=\frac{3 \mathrm{i}}{4 \pi} \int_{A_{01}^{\prime}} \mathrm{i} f \frac{\partial h}{\partial \varphi_{1}} \mathrm{~d} \varphi_{1} \wedge \mathrm{~d} \rho_{2}^{2} \wedge \mathrm{~d} \varphi_{2} \wedge \mathrm{~d} \rho_{3}^{2} \wedge \mathrm{~d} \varphi_{3} . \tag{3.14}
\end{equation*}
$$

The submanifold $A_{01}^{\prime}$ is oriented as a subset of $\partial B_{0}$ and the orientation of $B_{0}$ is the one defined by $\omega^{3}$, that is, by

$$
\mathrm{d} \rho_{1}^{2} \wedge \mathrm{~d} \varphi_{1} \wedge \mathrm{~d} \rho_{2}^{2} \wedge \mathrm{~d} \varphi_{2} \wedge \mathrm{~d} \rho_{3}^{2} \wedge \mathrm{~d} \varphi_{3}
$$

Since $\rho_{1}>\epsilon$ for the points of $B_{0}$, then $A_{01}^{\prime}$ is oriented by $-\mathrm{d} \varphi_{1} \wedge \mathrm{~d} \varphi_{2}^{2} \wedge \mathrm{~d} \varphi_{2} \wedge \mathrm{~d} \rho_{3}^{2} \wedge \mathrm{~d} \varphi_{3}$. On the other hand, the Hamiltonian function $f=-\kappa+O(\epsilon)$ on $A_{01}^{\prime}$. Then it follows from (3.14) together with (3.13)

$$
N_{01}^{\prime}=6 \pi^{2} \kappa \int_{0}^{\sigma / \pi} \mathrm{d} \rho_{3}^{2} \int_{0}^{\tau / \pi-\rho_{3}^{2}} \mathrm{~d} \rho_{2}^{2}+O(\epsilon)
$$

that is,

$$
\begin{equation*}
N_{01}^{\prime}=3 \kappa\left(\tau^{2}-\lambda^{2}\right)+O(\epsilon) \tag{3.15}
\end{equation*}
$$

The contributions $N_{02}^{\prime}, N_{03}^{\prime}, N_{04}^{\prime}, N_{05}^{\prime}$ to (3.12) can be calculated in a similar way. One obtains the following results up to addends of order $\epsilon$

$$
\begin{equation*}
N_{02}^{\prime}=N_{05}^{\prime}=-\left(\tau^{3}-\lambda^{3}\right)+3 \kappa\left(\tau^{2}-\lambda^{2}\right), \quad N_{03}^{\prime}=\tau^{2}(3 \kappa-\tau), \quad N_{04}^{\prime}=\lambda^{2}(3 \kappa-\lambda) . \tag{3.16}
\end{equation*}
$$

As $I_{\psi}$ is independent of $\epsilon$, it follows from (3.12), (3.15) and (3.16)

$$
\begin{equation*}
I_{\psi}=6 \kappa\left(2 \tau^{2}-\lambda^{2}\right)+\lambda^{3}-3 \tau^{3} \tag{3.17}
\end{equation*}
$$

On the other hand, straightforward calculations give

$$
\int_{M} \omega^{3}=\left(\tau^{3}-\lambda^{3}\right), \quad \text { and } \quad \int_{M} \pi \rho_{1}^{2} \omega^{3}=\frac{1}{4}\left(\tau^{4}-\lambda^{4}\right) .
$$

So

$$
\begin{equation*}
\kappa=\frac{1}{4}\left(\frac{\tau^{4}-\lambda^{4}}{\tau^{3}-\lambda^{3}}\right) . \tag{3.18}
\end{equation*}
$$

It follows from (3.17) and (3.18)

$$
\begin{equation*}
I_{\psi}=\frac{\lambda^{2}\left(-3 \tau^{4}+8 \tau^{3} \lambda-6 \tau^{2} \lambda^{2}+\lambda^{4}\right)}{2\left(\tau^{3}-\lambda^{3}\right)} \tag{3.19}
\end{equation*}
$$

Hence $I_{\psi}$ is a rational function of $\tau$ and $\lambda$. It is easy to check that its numerator does not vanish for $0<\lambda<\tau$. So we have proved the following proposition.

Proposition 8. If $\psi$ is the closed Hamiltonian isotopy defined in (3.7), then the characteristic number $I_{\psi} \neq 0$.
Proof of Corollary 2. By Proposition $8 I_{\psi} \neq 0$. As $I$ is a group homomorphism on $\pi_{1}(\operatorname{Ham}(M, \omega))$, then the class $\left[\psi^{l}\right] \in \pi_{1}(\operatorname{Ham}(M, \omega))$ does not vanish, for all $l \in \mathbb{Z} \backslash\{0\}$.

## 4. Hamiltonian group of toric manifolds

In this section we generalize the calculations carried out in Section 3 for the 6 -manifold one point blow up of $\mathbb{C} P^{3}$ to a general toric manifold. Now $(M, \omega)$ will denote the toric manifold defined by (1.3) and (1.4).

When $0 \neq z_{b} \in \mathbb{C}$, we write $z_{b}=\rho_{b} \mathrm{e}^{\mathrm{i} \theta_{b}}$. The standard symplectic form on $\mathbb{C}^{m}$ gives rise to the symplectic structure $\omega$ on $M$. On

$$
\left\{[z] \in M: z_{j} \neq 0 \text { for all } j\right\}
$$

$\omega$ can be written as in (3.5)

$$
\omega=\sum_{i=1}^{n} \mathrm{~d}\left(\frac{\rho_{a i}^{2}}{2}\right) \wedge \mathrm{d} \varphi_{a i},
$$

with $\varphi_{a i}$ a linear combination of the $\theta_{c}$ 's.
Given $0<\epsilon \ll 1$, we set

$$
\begin{aligned}
& B_{0}=\left\{[z] \in M:\left|z_{j}\right|>\epsilon \text { for all } j\right\} \\
& B_{k}=\left\{[z] \in M:\left|z_{k}\right|<2 \epsilon,\left|z_{j}\right|>\epsilon \text { for all } j \neq k\right\},
\end{aligned}
$$

as in Section 3. On $B_{0}$ we will consider the Darboux coordinates

$$
\left\{\frac{\rho_{a i}^{2}}{2}, \varphi_{a i}\right\}_{i=1, \ldots, n}
$$

Given $k \in\{1, \ldots, m\}$ we write $\omega$ in the form

$$
\omega=\mathrm{d}\left(\frac{\rho_{k}^{2}}{2}\right) \wedge \mathrm{d} \varphi_{k}+\sum_{i=1}^{n-1} \mathrm{~d}\left(\frac{\rho_{k i}^{2}}{2}\right) \wedge \mathrm{d} \varphi_{k i}
$$

where $\varphi_{k}$ and $\varphi_{k i}$ are linear combinations of the $\theta_{c}$ 's. Then we consider on $B_{k}$ the following Darboux coordinates

$$
\left\{x_{k}, y_{k}, \frac{\rho_{k i}^{2}}{2}, \varphi_{k i}\right\}_{i=1, \ldots, n-1}
$$

with $x_{k}, y_{k}$ defined by $x_{k}+\mathrm{i} y_{k}:=\rho_{k} \mathrm{e}^{\mathrm{i} \varphi_{k}}$, if $z_{k} \neq 0$ and $x_{k}=0=y_{k}$, if $z_{k}=0$.
We denote by $\psi_{t}$ the map

$$
\psi_{t}:[z] \in M \mapsto\left[z_{1} \mathrm{e}^{2 \pi \mathrm{i} t}, z_{2}, \ldots, z_{m}\right] \in M
$$

$\left\{\psi_{t}: t \in[0,1]\right\}$ is a loop in $\operatorname{Ham}(M)$. By repeating the arguments of Section 3 one obtains

$$
I_{\psi}=\sum_{k=1}^{m} N_{0 k}^{\prime}+O(\epsilon)
$$

where

$$
\begin{aligned}
& N_{0 k}^{\prime}=\frac{n \mathrm{i}}{2 \pi} \int_{A_{0 k}^{\prime}} f \mathrm{~d} \log r_{0 k} \wedge \omega^{n-1}, \\
& A_{0 k}^{\prime}=\left\{[z] \in M:\left|z_{k}\right|=\epsilon,\left|z_{j}\right|>\epsilon \text { for all } j \neq k\right\}
\end{aligned}
$$

and $f=\pi \rho_{1}^{2}-\kappa_{1}$, with

$$
\int_{M} \pi \rho_{1}^{2} \omega^{n}=\kappa_{1} \int_{M} \omega^{n}
$$

As in Section 3, on $A_{0 k}^{\prime}$ the exterior derivative $\mathrm{d} \log r_{0 k}=\mathrm{i} h^{\prime}\left(\varphi_{k}\right) \mathrm{d} \varphi_{k}$, where $h=h\left(\varphi_{k}\right)$ is a function such that $h(0)=2 \pi, h(2 \pi)=0$. Then

$$
N_{0 k}^{\prime}=-n \int_{\{[z]: z k=0\}} f \omega^{n-1}+O(\epsilon),
$$

where $\left\{[z] \in M: z_{k}=0\right\}$ is oriented by the restriction of $\omega$ to this submanifold. Since $I_{\psi}$ is independent of $\epsilon$, we obtain

$$
\begin{equation*}
I_{\psi}=-n \sum_{k=1}^{m}\left(\int_{\left\{[z]: z_{k}=0\right\}}\left(\pi \rho_{1}^{2}-\kappa_{1}\right) \omega^{n-1}\right) . \tag{4.1}
\end{equation*}
$$

For $j=1, \ldots, m$ we write

$$
\begin{equation*}
\alpha_{j}:=\sum_{k=1}^{m}\left(\int_{\{[z]: z k=0\}}\left(\pi \rho_{j}^{2}-\kappa_{j}\right) \omega^{n-1}\right), \tag{4.2}
\end{equation*}
$$

where $\kappa_{j}$ is defined by the condition

$$
\int_{M} \pi \rho_{j}^{2} \omega^{n}=\kappa_{j} \int_{M} \omega^{n}
$$

Proof of Theorem 3. Let us assume that $\alpha_{1}, \ldots, \alpha_{p}$ are linearly independent over $\mathbb{Z}$. For $j=1, \ldots, p$ we put

$$
{ }^{j} \psi_{t}:[z] \in M \mapsto\left[z_{1}, \ldots, z_{j} e^{2 \pi i t}, \ldots, z_{m}\right] \in M .
$$

Given $q=\left(q_{1}, \ldots, q_{p}\right) \in \mathbb{Z}^{p}$ we denote by $\psi^{q}$ the path product

$$
\left({ }^{1} \psi\right)^{q_{1}} \star \cdots \star\left({ }^{p} \psi\right)^{q_{p}} .
$$

Formula (4.1) together with the fact that $I$ is a group homomorphism give

$$
I_{\psi^{q}}=-n \sum_{i=1}^{p} q_{i} \alpha_{i}
$$

Analogously if $q^{\prime}=\left(q_{1}^{\prime}, \ldots, q_{p}^{\prime}\right) \in \mathbb{Z}^{p}$, then $I_{\psi q^{\prime}}=-n \sum_{i=1}^{p} q_{i}^{\prime} \alpha_{i}$. By the linear independence of $\alpha_{1}, \ldots, \alpha_{p}$ from $I_{\psi^{q^{\prime}}}=I_{\psi^{q}}$ it follows $q=q^{\prime}$. So $\psi^{q}$ is homotopic to $\psi^{q^{\prime}}$ iff $q=q^{\prime}$.

Example. We will check the above result calculating the family $\left\{\alpha_{j}\right\}$ defined in (4.2) in two particular cases: when the manifold is $\mathbb{C} P^{1}$ and when it is $\mathbb{C} P^{2}$.

For

$$
M=\mathbb{C} P^{1}=\left\{\left(z_{1}, z_{2}\right) \in \mathbb{C}^{2}:\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}=\tau / \pi\right\} / S^{1},
$$

we have

$$
\int_{M} \pi \rho_{1}^{2} \omega=\tau^{2} / 2, \quad \int_{M} \omega=\tau
$$

Thus $\kappa_{1}=\tau / 2$ and $\alpha_{1}=-\kappa_{1}+\tau-\kappa_{1}=0$. Similarly $\alpha_{2}=0$. In this case the number $p$ in Theorem 3 is 0 . This is compatible with the fact that $\pi_{1}\left(\operatorname{Ham}\left(\mathbb{C} P^{1}\right)\right)=\mathbb{Z} / 2 \mathbb{Z}$.

For

$$
M=\mathbb{C} P^{2}=\left\{\left(z_{1}, z_{2}, z_{3}\right) \in \mathbb{C}^{3}:\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}=\tau / \pi\right\} / S^{1},
$$

we have the following values for the integrals involved in the definition of $\alpha_{1}$

$$
\int_{M} \omega^{2}=\tau^{2}, \quad \int_{M} \pi \rho_{1}^{2} \omega^{2}=\tau^{3} / 3
$$

So $\kappa_{1}=\tau / 3$. Moreover for $k \in\{1,2,3\}$

$$
\int_{\left\{[z]: z_{k}=0\right\}} \omega=\tau .
$$

On the other hand, for $k=2,3$

$$
\int_{\{[z]: ; z=0\}} \pi \rho_{1}^{2} \omega=\tau^{2} / 2
$$

So $\alpha_{1}=-\kappa_{1} \tau+\left(\tau^{2} / 2-\kappa_{1} \tau\right)+\left(\tau^{2} / 2-\kappa_{1} \tau\right)=0$. Analogously $\alpha_{2}=\alpha_{3}=0$, so $p$ in Theorem 3 is also 0 . This result is consistent with the finiteness of $\pi_{1}\left(\operatorname{Ham}\left(\mathbb{C} P^{2}\right)\right)$, for $\operatorname{Ham}\left(\mathbb{C} P^{2}\right)$ has the homotopy type of $P U(3)$ [9].

Remark. On the manifold $M$ one point blow up of $\mathbb{C} P^{3}$, defined by (3.1) and (3.2), one can consider the loop $\tilde{\psi}$ defined by

$$
\begin{equation*}
\tilde{\psi}_{t}[z]=\left[z_{1}, z_{2}, z_{3} \mathrm{e}^{2 \pi \mathrm{i} t}, z_{4}, z_{5}\right] . \tag{4.3}
\end{equation*}
$$

A similar calculation to the one carried out in the proof of (3.19) shows that $I_{\tilde{\psi}}=-3 I_{\psi}$.
In the definition of $M$ the variables $z_{1}, z_{2}, z_{5}$ play the same role. However we can consider the following $S^{1}$-action on $M$

$$
\begin{equation*}
\hat{\psi}_{t}[z]=\left[z_{1}, z_{2}, z_{3}, z_{4} \mathrm{e}^{2 \pi i t}, z_{5}\right], \tag{4.4}
\end{equation*}
$$

and it turns out that $I_{\hat{\psi}}=3 I_{\psi}$. Thus Theorem 3 guaranties that only $\mathbb{Z}$ is contained in $\pi_{1}(\operatorname{Ham}(M))$.
Let $(M, \omega)$ be the toric manifold determined by the Delzant polytope $\Delta \subset \mathfrak{t}^{*}$, where $T$ is an $n$-dimensional torus. Next we give a formula for the value of $I$ on the Hamiltonian loops generated by the effective action of $T$ on $M$, in which are involved geometrical magnitudes relative to $\Delta$ and the generator of the loop.

By $\mu: M \rightarrow \mathfrak{t}^{*}$ is denoted the moment map for the $T$-action. Let $\mathbf{b}$ be an element of the integer lattice of $\mathfrak{t}$, and let $\psi_{\mathbf{b}}$ the $S^{1}$-action determined by $\mathbf{b}$. The corresponding normalized Hamiltonian function is $f=\langle\mu, \mathbf{b}\rangle-\kappa$, with

$$
\int_{M}\langle\mu, \mathbf{b}\rangle \omega^{n}=\kappa \int_{M} \omega^{n}
$$

Since

$$
\int_{M} \mu \omega^{n}=\operatorname{Cm}(\Delta) \int_{M} \omega^{n}
$$

where $\mathrm{Cm}(\Delta)$ is the center of mass of $\Delta$, it follows $\kappa=\langle\mathrm{Cm}(\Delta), \mathbf{b}\rangle$.
According to (4.1)

$$
I_{\psi_{\mathbf{b}}}=-n \sum_{k=1}^{m} \int_{D_{k}}(\langle\mu, \mathbf{b}\rangle-\langle\mathrm{Cm}(\Delta), \mathbf{b}\rangle) \omega^{n-1}
$$

where $D_{k}:=\mu^{-1}\left(F_{k}\right)$, and $F_{1}, \ldots, F_{m}$ are the facets of $\Delta$.
We define $\mathrm{C} \mathrm{m}\left(D_{k}\right)$ by the relation

$$
\operatorname{Cm}\left(D_{k}\right) \int_{D_{k}} \omega^{n-1}=\int_{D_{k}} \mu \omega^{n-1}
$$

and

$$
\operatorname{Vol}\left(D_{k}\right):=\frac{1}{(n-1)!} \frac{1}{(2 \pi)^{n-1}} \int_{D_{k}} \omega^{n-1}
$$

Then

$$
\begin{equation*}
I_{\psi_{\mathbf{b}}}=n!(2 \pi)^{n-1} \sum_{k=1}^{m}\left\langle\mathrm{Cm}(\Delta)-\mathrm{Cm}\left(D_{k}\right), \mathbf{b}\right\rangle \operatorname{Vol}\left(D_{k}\right) \tag{4.5}
\end{equation*}
$$

Thus we have the following proposition.

Proposition 9. Let $(M, \omega)$ be the toric manifold associated to the polytope $\Delta$. If there is $\mathbf{b}$ in the integer lattice of $\mathfrak{t}$ such that

$$
\sum_{k=1}^{m}\left\langle\mathrm{Cm}(\Delta)-\operatorname{Cm}\left(D_{k}\right), \mathbf{b}\right\rangle \operatorname{Vol}\left(D_{k}\right) \neq 0
$$

then $\mathbf{b}$ generates an element of infinite order in the group $\pi_{1}(\operatorname{Ham}(M, \omega))$.

## 5. Hamiltonian $G$-actions

Let $G$ be a compact Lie group and $\phi: G \rightarrow \operatorname{Ham}(M, \omega)$ a Hamiltonian $G$-action on $M$. The group homomorphism $\phi$ induces a map

$$
\Phi: B G \rightarrow B \operatorname{Ham}(M, \omega)
$$

between the corresponding classifying spaces.
On the other hand, one has the universal bundle with fibre $M$

where $E \operatorname{Ham}(M) \rightarrow B \operatorname{Ham}(M)$ is the universal principal bundle of the group $H:=\operatorname{Ham}(M, \omega)$.
The pullback $\Phi^{-1}\left(M_{H}\right)$ of $M_{H}$ by $\Phi$ is a bundle on $B G$ which can be identified with $p: M_{\phi}:=E G \times_{G} M \rightarrow$ $B G$. Thus we have the following commutative diagram


There exists a unique class $\mathbf{c} \in H^{2}\left(M_{H}, \mathbb{R}\right)$ [14] called the coupling, such that $\mathbf{c}$ extends the fiberwise class $[\omega]$ and $\pi_{H_{*}} \mathbf{c}^{n+1}=0$ (where $\pi_{H_{*}}$ is the fiber integration). We put $c_{\phi}$ for the pullback of $\mathbf{c}$ by $\Phi^{\prime}$; that is, $c_{\phi}=\Phi^{*}(\mathbf{c}) \in H^{2}\left(M_{\phi}, \mathbb{R}\right)$. Since

$$
p_{*}\left(c_{\phi}^{n+1}\right)=\Phi^{*}\left(\pi_{H_{*}} \mathbf{c}^{n+1}\right)=0,
$$

$c_{\phi}$ is the coupling class of the Hamiltonian fibration $M_{\phi} \rightarrow B G$ [21].
We can also consider the vector bundle

$$
(T M)_{\phi}:=E G \times_{G} T M \rightarrow M_{\phi} .
$$

The first Chern class $c_{1}\left((T M)_{\phi}\right)$ is the $G$-equivariant first Chern class of $T M$, and it will be denoted by $c_{1}^{\phi}$.
By $\operatorname{Hom}(G, \operatorname{Ham}(M))$ is denoted the set of all Lie group homomorphisms $\phi$ from $G$ to $\operatorname{Ham}(M, \omega)$. In $\operatorname{Hom}(G, \operatorname{Ham}(M))$ one defines the following equivalence relation:
$\phi \simeq \tilde{\phi}$ iff there is a continuous family $\left\{\phi^{s}: G \rightarrow \operatorname{Ham}(M)\right\}_{s \in[0,1]}$ of Lie group homomorphisms, such that $\phi^{0}=\phi$ and $\phi^{1}=\tilde{\phi}$; that is, iff $\phi$ and $\tilde{\phi}$ are homotopic by a family of group homomorphisms. We denote by $[G, \operatorname{Ham}(M)]_{g h}$ the corresponding quotient set. This space is just a set of connected components of the space of homomorphisms from $G$ to $\operatorname{Ham}(\underset{\sim}{M}, \omega)$.

If $\phi \simeq \tilde{\phi}$, then the bundles $\tilde{\Phi}^{-1}\left(M_{H}\right)$ and $\Phi^{-1}\left(M_{H}\right)$ are isomorphic. Moreover the isomorphism $M_{\phi} \rightarrow M_{\tilde{\phi}}$ applies $c_{\tilde{\phi}}$ in $c_{\phi}$ and $c_{1}^{\tilde{\phi}}$ in $c_{1}^{\phi}$.

For $j=0,1, \ldots, n$ we put

$$
\beta_{j}(\phi):=\left(c_{1}^{\phi}\right)^{j}\left(c_{\phi}\right)^{n-j} \in H^{2 n}\left(M_{\phi}, \mathbb{R}\right)
$$

We write $R_{j}(\phi):=p_{*}\left(\beta_{j}(\phi)\right) \in H^{0}(B G)$. By the localization formula in $G$-equivariant cohomology $[5,13]$

$$
\begin{equation*}
R_{i}(\phi)=\sum_{Z} p_{*}^{Z}\left(\frac{\left.\beta\right|_{Z}}{e_{Z}}\right) \tag{5.1}
\end{equation*}
$$

where $Z$ varies in the set of connected components of the fixed point set, $p_{*}^{Z}: H_{G}(Z) \rightarrow H(B G)$ is the fiber integration on $Z$, and $e_{Z}$ is the equivariant Euler class of the normal bundle to $Z$ in $M$.

From the preceding arguments it follows the following theorem.
Theorem 10. Given $\phi$ and $\tilde{\phi}$ two Hamiltonian $G$-actions on $M$, if there are $j \in\{0,1, \ldots, n\}$ and $X \in \mathfrak{g}$ such that $R_{j}(\phi)(X) \neq R_{j}(\tilde{\phi})(X)$, then $[\phi] \neq[\tilde{\phi}] \in[G, \operatorname{Ham}(M)]_{g h}$.

If $\hat{\omega} \in H^{2}\left(M_{\phi}, \mathbb{R}\right)$ is an element which restricts to the class of the symplectic form on the fiber in the fibration $p: M_{\phi} \rightarrow B G$, then

$$
c_{\phi}=\hat{\omega}-\frac{1}{k} p^{*}\left(p_{*}\left(\hat{\omega}^{n+1}\right)\right),
$$

where the constant $k=(n+1) \int_{M} \omega^{n}$ (see [14]). In particular, if $G=U(1)$ we denote by $f$ the normalized Hamiltonian function; that is, $\iota_{Y} \omega=-\mathrm{d} f$ and $\int_{M} f \omega^{n}=0$, where $Y$ the vector field on $M$ generated by $\phi$. Then $c_{\phi}$ is the class in $H^{2}\left(M_{\phi}\right)$ defined by the $U(1)$-equivariant 2-form $\omega+f u$, where $u$ is a coordinate on the Lie algebra $\mathfrak{u}(1)$ dual of a fixed base $X$ of $\mathfrak{u}(1)$ (see $[16,13]$ ).

When $G=U(1)$ a representative of $c_{1}^{\phi}(\operatorname{det}(T M))$ can be constructed following [3] or [5]. Let $s$ be a local section of $\operatorname{det}(T M)$ over the open $V$. The infinitesimal action of $X$ on the section $s$ is the section $X s$ defined in (1.11). $X s$ is a section which can be written as the product $L \cdot s$, of a function $L$ on $V$ and $s$. If $\alpha$ is the form relative to $s$ of an equivariant connection on $\operatorname{det}(T M)$, and $X_{M}$ is the Hamiltonian vector field on $M$ defined in (1.9), then

$$
\begin{equation*}
\frac{-1}{2 \pi \mathrm{i}}\left(\mathrm{~d} \alpha+\left(L-\iota_{X_{M}} \alpha\right) u\right), \tag{5.2}
\end{equation*}
$$

is a representative of $c_{1}^{\phi}(\operatorname{det}(T M))$ on $V$. So a representative of $\beta_{1}=c_{1}^{\phi}(T M) c_{\phi}^{n-1}$ on $V$ is

$$
\begin{equation*}
\frac{-1}{2 \pi \mathrm{i}}\left(\mathrm{~d} \alpha+\left(L-\iota_{X_{M}} \alpha\right) u\right) \wedge(\omega+f u)^{n-1} . \tag{5.3}
\end{equation*}
$$

On the other hand, if $G=U(1)$ and $\mu: M \rightarrow \mathfrak{u}(1)^{*}$ is the normalized moment map, $\mathcal{R}_{\phi}(Z)$ defined in (1.6) is equal to

$$
\begin{equation*}
\mathcal{R}_{\phi}(Z)=\left(\frac{1}{2 \pi \mathrm{i}}\right)^{n} \int_{M} \mathrm{e}^{\mathrm{i} c_{\phi}(Z)} \tag{5.4}
\end{equation*}
$$

for any $Z \in \mathfrak{g}$. So Theorem 10 is applicable to $\mathcal{R}$.

### 5.1. Flag manifolds

Let $\eta \in \mathfrak{g}^{*}$ be a regular element; that is, the stabilizer $G_{\eta}$ of $\eta$ for the coadjoint action of $G$ is a maximal torus $T$. By $\mathcal{O}$ is denoted the coadjoint orbit of $\eta$, endowed with the Kirillov symplectic structure $\omega$. The $G$-action on $\mathcal{O}$ is Hamiltonian and the inclusion map $\mu: \mathcal{O} \rightarrow \mathfrak{g}^{*}$ is a moment map for this action. The Fourier transform of the orbit $\mathcal{O}$ is the function $F$ defined on $\mathfrak{g}$ by (see [5])

$$
\begin{equation*}
F(X)=\left(\frac{1}{2 \pi \mathrm{i}}\right)^{n} \int_{\mathcal{O}} \mathrm{e}^{\mathrm{i}(\mu(X)+\omega)}, \tag{5.5}
\end{equation*}
$$

where $X \in \mathfrak{g}$ and $n=(\operatorname{dim} \mathcal{O}) / 2$.
Let $Y$ be a vector of $\mathfrak{g}$, by $\phi_{t}^{Y}$ we denote the isotopy defined in (1.10). If $\left\{\phi_{t}^{Y}\right\}_{t \in[0,1]}$ is a closed curve in $\operatorname{Ham}(\mathcal{O})$, we have a Hamiltonian circle action $\phi^{Y}: U(1) \rightarrow \operatorname{Ham}(\mathcal{O})$ and $\mu(Y)$ is a Hamiltonian function for this $S^{1}$-action. If

$$
\begin{equation*}
\kappa:=\left(\int_{\mathcal{O}} \mu(Y) \omega^{n}\right)\left(\int_{\mathcal{O}} \omega^{n}\right)^{-1} \tag{5.6}
\end{equation*}
$$

then $f=\mu(Y)-\kappa$ is the normalized Hamiltonian which generates the $U(1)$-action.

On the other hand, one deduces from (5.5)

$$
\begin{equation*}
\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} F(t Y)=\left(\frac{1}{2 \pi}\right)^{n} \frac{\mathrm{i}}{n!} \int_{\mathcal{O}} \mu(Y) \omega^{n} . \tag{5.7}
\end{equation*}
$$

It follows from (5.6) and (5.7) the following formula for the constant $\kappa$

$$
\begin{equation*}
\kappa=\left.\frac{-\mathrm{i}}{\operatorname{Vol}(\mathcal{O})} \frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} F(t Y), \tag{5.8}
\end{equation*}
$$

where the symplectic volume is

$$
\operatorname{Vol}(\mathcal{O})=\frac{1}{(2 \pi)^{n}} \frac{1}{n!} \int_{\mathcal{O}} \omega^{n} .
$$

According to (5.4) and (5.8) we have the following proposition.
Proposition 11. Given $Y \in \mathfrak{g}$, if $\phi^{Y}$ is a loop in $\operatorname{Ham}(\mathcal{O})$, then

$$
\begin{equation*}
\mathcal{R}_{\phi^{Y}}(Y)=\exp \left(-\left.\frac{\mathrm{d}}{\mathrm{~d} t}\right|_{t=0} \log F(t Y)\right) F(Y) \tag{5.9}
\end{equation*}
$$

Let $W$ be the Weyl group determined by the torus $T$ and $X$ an regular element of $\mathfrak{t}$. The Harish-Chandra theorem gives a formula for $F(X)$ in terms of roots of $\mathfrak{t}$ (see [5])

$$
\begin{equation*}
F(X)=\prod_{\eta(i \dot{\alpha})>0}(\alpha(X))^{-1} \sum_{w \in W} \epsilon(w) \mathrm{e}^{\mathrm{i}(w \cdot \eta)(X)}, \tag{5.10}
\end{equation*}
$$

where $\epsilon(w)$ is the signature of the permutation $w \in W$. From (5.9) and (5.10) one deduces that $\mathcal{R}_{\phi^{Y}}(Y)$ can be calculated using the root structure defined by the pair ( $G, T$ ).

Example. Let $G=S U(2)$ and $\eta \in \mathfrak{s} u(2)^{*}$ defined by

The stabilizer of $\eta$ is $T=U(1)$, the coadjoint orbit $\mathcal{O}$ is $\mathbb{C} P^{1}$ and the corresponding symplectic form is $\omega_{\text {area. }}$. The Weyl group $W=N(T) / T$, with $N(T)$ the normalizer of $T$ in $S U(2)$, consists of the class of id and the class of

$$
\left(\begin{array}{cc}
0 & 1  \tag{5.11}\\
-1 & 0
\end{array}\right) .
$$

If $\alpha_{1}=\epsilon_{1}-\epsilon_{2}$ is the usual base of roots, then $\eta\left(\mathrm{i} \check{\alpha}_{1}\right)=1$. In this case the product in (5.10) has only one factor and the sum of two addends.

Let $Y=\operatorname{diag}(\pi \mathrm{i},-\pi \mathrm{i})$, then the vector $C_{t}$ in (2.3) for $n=k=1$ is equal to $t Y$. Thus $\phi^{Y}$ is the loop in $\operatorname{Ham}(\mathcal{O})$ denoted by ${ }_{1} \psi$ in Section 2. This loop defines the only nontrivial class of $\pi_{1}(\operatorname{Ham}(\mathcal{O})$ ) (see the paragraph before Theorem 5).

If $w$ is the element of $W$ defined by (5.11), then $w Y=\operatorname{diag}(-\pi \mathrm{i}, \pi \mathrm{i})$, and $\eta(w Y)=-\pi$. Furthermore $\alpha_{1}(Y)=2 \pi \mathrm{i}$. It follows from (5.10)

$$
F(Y)=\frac{1}{2 \pi \mathrm{i}}\left(\mathrm{e}^{\pi \mathrm{i}}-\mathrm{e}^{-\pi \mathrm{i}}\right)=0
$$

By (5.9) $\mathcal{R}_{\phi^{Y}}(Y)=0$.
In general, if $b \neq 0$ then for $Z=b Y$,

$$
F(Z)=\frac{\sin b \pi}{b \pi}
$$

The loop determined by $2 Y$ defines the trivial class in $\pi_{1}(\operatorname{Ham}(\mathcal{O}))$, and $\mathcal{R}_{\phi^{2 Y}}(2 Y)=0$.

On the other hand,

$$
\begin{equation*}
\mathcal{R}_{\phi^{0}}(X)=\operatorname{Vol}(\mathcal{O}), \tag{5.12}
\end{equation*}
$$

for any $0 \neq X \in \mathfrak{g}$. It follows from Theorem 10 and (5.12) the following proposition.
Proposition 12. The circle action on $\left(\mathbb{C} P^{1}, \omega_{\text {area }}\right)$ defined by $\operatorname{diag}(2 \pi \mathrm{i},-2 \pi \mathrm{i}) \in \mathfrak{s u}(2)$ and the trivial one determine distinct elements in $\left[U(1), \operatorname{Ham}\left(\mathbb{C} P^{1}\right)\right]_{g h}$.

Proposition 12 gives an example of a pair of Hamiltonian circle actions on $\left(\mathbb{C} P^{1}, \omega_{\text {area }}\right)$ which define the same element in $\pi_{1}(\mathrm{Ham})$ but they are not homotopic by a family of circle actions.

### 5.2. Hirzebruch surfaces

Next we determine the value $R_{\phi}:=R_{1}(\phi)$ for three Hamiltonian actions on a Hirzebruch surface. We will define these actions using the fact that such a surface is a submanifold of $\mathbb{C} P^{1} \times \mathbb{C} P^{2}$.

Given 3 numbers $k, \tau, \sigma$, with $k \in \mathbb{Z}_{>0}, \tau, \sigma \in \mathbb{R}_{>0}$ and $k \sigma<\tau$, the triple $(k, \tau, \sigma)$ determine a Hirzebruch surface $M$ (see [4]). This surface is the quotient

$$
\left\{z \in \mathbb{C}^{4}: k\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\left|z_{4}\right|^{2}=\tau / \pi,\left|z_{1}\right|^{2}+\left|z_{3}\right|^{2}=\sigma / \pi\right\} / \mathbb{T}
$$

where the equivalence defined by $\mathbb{T}=\left(S^{1}\right)^{2}$ is given by

$$
(a, b) \cdot\left(z_{1}, z_{2}, z_{3}, z_{4}\right)=\left(a^{k} b z_{1}, a z_{2}, b z_{3}, a z_{4}\right),
$$

for $(a, b) \in\left(S^{1}\right)^{2}$.
The map

$$
\left[z_{1}, z_{2}, z_{3}, z_{4}\right] \mapsto\left(\left[z_{2}: z_{4}\right],\left[z_{2}^{k} z_{3}: z_{4}^{k} z_{3}: z_{1}\right]\right)
$$

allows us to represent $M$ as a submanifold of $\mathbb{C} P^{1} \times \mathbb{C} P^{2}$. On the other hand, the usual symplectic structures on $\mathbb{C} P^{1}$ and $\mathbb{C} P^{2}$ induce a symplectic form $\omega$ on $M$, and the following $\left(S^{1}\right)^{2}$-action on $\mathbb{C} P^{1} \times \mathbb{C} P^{2}$

$$
(a, b)\left(\left[u_{0}: u_{1}\right],\left[x_{0}: x_{1}: x_{2}\right]\right)=\left(\left[a u_{0}: u_{1}\right],\left[a^{k} x_{0}: x_{1}: b x_{2}\right]\right)
$$

gives rise to a toric structure on $M$. The Delzant polytope associated to $(M, \omega)$ is the trapezoid in $\left(\mathbb{R}^{2}\right)^{*}$ whose not oblique edges have the lengths $\tau, \sigma$, and $\lambda:=\tau-k \sigma$ (see [10]). Moreover $\lambda$ is the value that the symplectic form $\omega$ takes on $\left\{[z] \in M: z_{3}=0\right\}$, the exceptional divisor of $M$, when $k=1$. And $\omega$ takes the value $\sigma$ on the class of the fibre in the fibration $M \rightarrow \mathbb{C} P^{1}$.

Let $\phi_{t}$ be the diffeomorphism of $M$ defined by

$$
\begin{equation*}
\phi_{t}\left[z_{1}, z_{2}, z_{3}, z_{4}\right]=\left[z_{1} \mathrm{e}^{2 \pi \mathrm{i} t}, z_{2}, z_{3}, z_{4}\right] . \tag{5.13}
\end{equation*}
$$

$\phi=\left\{\phi_{t}: t \in[0,1]\right\}$ is a loop of Hamiltonian symplectomorphisms of $(M, \omega)$. The fixed point set is $Z=\{[z] \in M$ : $\left.z_{1}=0\right\}$; that is, $Z \simeq \mathbb{C} P^{1}$ is the section at infinity of $M \rightarrow \mathbb{C} P^{1}$ (see [4]).

On $M$ we can consider the covering

$$
\begin{array}{ll}
U_{1}=\left\{[z] \in M: z_{3} \neq 0 \neq z_{4}\right\}, & U_{2}=\left\{[z] \in M: z_{1} \neq 0 \neq z_{4}\right\} \\
U_{3}=\left\{[z] \in M: z_{1} \neq 0 \neq z_{2}\right\}, & U_{4}=\left\{[z] \in M: z_{2} \neq 0 \neq z_{3}\right\} .
\end{array}
$$

So $Z \cap U_{j}=\emptyset$, for $j=2$, 3. On $U_{4}$ one defines the complex coordinates

$$
w_{0}:=\frac{z_{4}}{z_{2}}, \quad w_{0}^{\prime}:=\frac{z_{1}}{z_{3} z_{2}^{k}} .
$$

In these coordinates

$$
\begin{equation*}
\phi_{t}\left(w_{0}, w_{0}^{\prime}\right)=\left(w_{0}, w_{0}^{\prime} \mathrm{e}^{2 \pi \mathrm{i} t}\right) . \tag{5.14}
\end{equation*}
$$

On $U_{1}$ we introduce the complex coordinates

$$
\begin{equation*}
w_{1}:=\frac{z_{2}}{z_{4}}, \quad w_{1}^{\prime}:=\frac{z_{1}}{z_{3} z_{4}^{k}} . \tag{5.15}
\end{equation*}
$$

Thus on $U_{1} \cap U_{4}$ one has the following relation

$$
\begin{equation*}
\frac{\partial}{\partial w_{1}} \wedge \frac{\partial}{\partial w_{1}^{\prime}}=-w_{0}^{k+2} \frac{\partial}{\partial w_{0}} \wedge \frac{\partial}{\partial w_{0}^{\prime}} \tag{5.16}
\end{equation*}
$$

between the sections of $\operatorname{det}(T M)$.
On $Z \cap U_{4}$ we have the complex coordinate $w_{0}$ and on $Z \cap U_{1}$ the coordinate $w_{1}$, with $w_{1}=w_{0}^{-1}$. By (5.16) the bundle $\left.\operatorname{det}(T M)\right|_{Z}$ is the one whose first Chern class is $(k+2)$; that is, $\left.\operatorname{det}(T M)\right|_{Z}=\mathcal{O}(k+2)$.

Let us consider the local section

$$
s:=\frac{\partial}{\partial w_{0}} \wedge \frac{\partial}{\partial w_{0}^{\prime}}
$$

of $\operatorname{det}(T M)$. We need to determine the corresponding function $L$ which appears in (5.2). From (5.14) it follows

$$
\begin{equation*}
\left(\phi_{t}\right)_{*}\left(\frac{\partial}{\partial w_{0}}\right)=\frac{\partial}{\partial w_{0}}, \quad\left(\phi_{t}\right)_{*}\left(\frac{\partial}{\partial w_{0}^{\prime}}\right)=\mathrm{e}^{2 \pi \mathrm{it}} \frac{\partial}{\partial w_{0}^{\prime}}, \tag{5.17}
\end{equation*}
$$

then

$$
\left(\phi_{t}\right)_{*}(s)=\mathrm{e}^{2 \pi \mathrm{i} t} s
$$

Thus the above function $L$ is the constant $2 \pi \mathrm{i}$.
On the other hand, the Hamiltonian vector field $X_{M}$ which corresponds to $X=2 \pi \mathfrak{i} \in \mathfrak{u}(1)$ is

$$
X_{M}=-2 \pi \mathrm{i} w_{0}^{\prime} \frac{\partial}{\partial w_{0}^{\prime}}
$$

On $Z^{\prime}:=Z \backslash\left(\left\{w_{0}=0\right\} \cup\left\{w_{1}=0\right\}\right)$ the class $\beta=c_{1}^{\phi}(T M) c_{\phi}$ is represented by the equivariant form

$$
\begin{equation*}
\left(\left.\delta\right|_{Z^{\prime}}-\frac{1}{2 \pi \mathrm{i}}(2 \pi \mathrm{i}-0) u\right)\left(\left.\omega\right|_{Z^{\prime}}+\left.f\right|_{Z^{\prime}} u\right), \tag{5.18}
\end{equation*}
$$

where $\delta$ is a 2 -form representing the ordinary first Chern class of $\operatorname{det}(T M)$.
The normalized Hamiltonian function is $f=\pi\left|z_{1}\right|^{2}-\kappa$, with $\kappa \in \mathbb{R}$. So $f_{\mid Z}=-\kappa$. The constant $\kappa$ is fixed by the normalization condition. An easy calculation gives

$$
\begin{equation*}
\kappa=\frac{\sigma}{3}\left(\frac{3 \lambda+k \sigma}{2 \lambda+k \sigma}\right) . \tag{5.19}
\end{equation*}
$$

Next we calculate the equivariant Euler class $e_{Z}$ of the normal bundle $N_{Z}$ to $Z$ in $M$. Let $q$ be a point of $Z^{\prime}$, then

$$
T_{q} M=\mathbb{C} \frac{\partial}{\partial w_{0}^{\prime}} \oplus T_{q} Z
$$

Since $\frac{\partial}{\partial w_{0}^{\prime}}$ and $\frac{\partial}{\partial w_{1}^{\prime}}$ are sections of $N_{Z}$ on $Z \cap U_{4}$ and $Z \cap U_{1}$ respectively, such that

$$
\frac{\partial}{\partial w_{1}^{\prime}}=w_{0}^{k} \frac{\partial}{\partial w_{0}^{\prime}},
$$

we have $N_{Z}=\mathcal{O}(k)$.

We put $w_{0}^{\prime}=x+\mathrm{i} y$ and on $N_{Z}$ we consider the orientation defined by $\frac{\partial}{\partial w_{0}^{\prime}}$. The vector field $X_{M}$ at $q=(x, y)$ is $X_{M}=2 \pi\left(y \frac{\partial}{\partial x}-x \frac{\partial}{\partial y}\right)$. So

$$
\left[X_{M}, \frac{\partial}{\partial x}\right]=2 \pi \frac{\partial}{\partial y}, \quad\left[X_{M}, \frac{\partial}{\partial y}\right]=-2 \pi \frac{\partial}{\partial x} .
$$

Then $\operatorname{det}^{1 / 2}\left(\left[X_{M},\right]\right)=2 \pi$. If $c_{1}(\mathcal{O}(k))$ is represented in $Z^{\prime}$ by the 2 -form $\chi$, then the equivariant Euler class $e_{Z}$ is represented by (see [5])

$$
\chi+(-2 \pi) \operatorname{det}^{-1 / 2}\left(\left[X_{M},\right]\right) u=\chi-u .
$$

It follows from (5.1) and (5.18) that

$$
\begin{equation*}
R_{\phi}=\int_{Z^{\prime}} \frac{(\delta-u)(\omega-\kappa u)}{\chi-u}=\frac{-1}{u} \int_{Z^{\prime}}(\delta-u)(\omega-\kappa u)(1+\chi / u)=2 \kappa+\omega(Z) . \tag{5.20}
\end{equation*}
$$

A straightforward calculation gives $\int_{Z} \omega=\lambda+k \sigma$. It follows from this value together with (5.19) and (5.20)

$$
\begin{equation*}
R_{\phi}=\frac{6 \lambda^{2}+(9 k+6) \lambda \sigma+\left(2 k+3 k^{2}\right) \sigma^{2}}{6 \lambda+3 k \sigma} . \tag{5.21}
\end{equation*}
$$

Given $0 \neq r \in \mathbb{Z}$ we can consider the loop $\xi$ defined by

$$
\xi_{t}[z]=\left[z_{1} \mathrm{e}^{2 \pi r t \mathrm{i}}, z_{2}, z_{3}, z_{4}\right] .
$$

The fixed point set for this $U(1)$-action set is $Z$ as well. The corresponding Hamiltonian action is $r f$. In this case the respective function $L$ is $2 \pi r \mathrm{i}$, and now the equivariant Euler class of $N_{Z}$ is $e_{Z}=c_{1}(\mathcal{O}(k))-r u$. Hence

$$
\begin{equation*}
R_{\xi}=\frac{-1}{r u} \int_{Z^{\prime}}(\delta-r u)(\omega-r \kappa u)(1+\chi /(r u))=R_{\phi} \tag{5.22}
\end{equation*}
$$

Next we shall determine $R_{\tilde{\phi}}$, where the $U(1)$-action $\tilde{\phi}_{t}$ is defined by

$$
\begin{equation*}
\tilde{\phi}_{t}[z]=\left[z_{1}, z_{2} \mathrm{e}^{2 \pi t \mathrm{i}}, z_{3}, z_{4}\right] . \tag{5.23}
\end{equation*}
$$

Now the fixed point set is $\tilde{Z}=\left\{[z] \in M: z_{2}=0\right\}$, it is the fibre over $[0: 1]$ of the fibration $M \rightarrow \mathbb{C} P^{1}$ and can be identified with

$$
\mathbb{C} P^{1} \simeq\left\{\left([0: 1],\left[0: z_{3}: z_{1}\right]\right)\right\} \subset \mathbb{C} P^{1} \times \mathbb{C} P^{2}
$$

The normalized Hamiltonian function is $\tilde{f}=\pi\left|z_{2}\right|^{2}-\tilde{\kappa}$, with

$$
\begin{equation*}
\tilde{\kappa}=\frac{3 \lambda^{2}+3 k \lambda \sigma+k^{2} \sigma^{2}}{6 \lambda+3 k \sigma} . \tag{5.24}
\end{equation*}
$$

A calculation similar to the preceding one shows that

$$
\begin{equation*}
R_{\tilde{\psi}}=2 \tilde{\kappa}+\omega(\tilde{Z}) . \tag{5.25}
\end{equation*}
$$

Since $\int_{\tilde{Z}} \omega=\sigma$, it follows from (5.25) together with (5.24) that

$$
\begin{equation*}
R_{\tilde{\psi}}=\frac{6 \lambda^{2}+(6 k+6) \lambda \sigma+\left(3 k+2 k^{2}\right) \sigma^{2}}{6 \lambda+3 k \sigma} . \tag{5.26}
\end{equation*}
$$

One can consider the Hamiltonian loop $\hat{\phi}_{t}$ defined by

$$
\begin{equation*}
\hat{\phi}_{t}[z]=\left[z_{1}, z_{2}, z_{3} \mathrm{e}^{2 \pi \mathrm{i} t}, z_{4}\right] . \tag{5.27}
\end{equation*}
$$

The corresponding fixed point set is $\hat{Z}=\left\{[z]: z_{3}=0\right\}$. The normalized Hamiltonian function $\hat{f}$ is $\hat{f}=\pi\left|z_{3}\right|^{2}-\hat{\kappa}$, with

$$
\begin{equation*}
\hat{\kappa}=\frac{3 \lambda \sigma+2 k \sigma^{2}}{6 \lambda+3 k \sigma} . \tag{5.28}
\end{equation*}
$$

It is easy to prove that

$$
\begin{equation*}
R_{\hat{\phi}}=2 \hat{\kappa}+\omega(\hat{Z}) \tag{5.29}
\end{equation*}
$$

and $\omega(\hat{Z})=\lambda$.
We can state the following theorem.
Theorem 13. If $\phi_{t}, \tilde{\phi}_{t}$ and $\hat{\phi}_{t}$ are the loops in $\operatorname{Ham}(M)$ defined by (5.13), (5.23) and (5.27) respectively, then

$$
R_{\phi}=2 \kappa+\omega(Z), \quad R_{\tilde{\phi}}=2 \tilde{\kappa}+\omega(\tilde{Z}), \quad R_{\hat{\phi}}=2 \hat{\kappa}+\omega(\hat{Z}),
$$

where $Z, \tilde{Z}$, and $\hat{Z}$ are the respective fixed point sets and the constants $\kappa, \tilde{\kappa}$, and $\hat{\kappa}$ are given by (5.19), (5.24) and (5.28) respectively.

From (5.21) and (5.26) it follows

$$
\begin{equation*}
R_{\phi}-R_{\tilde{\phi}}=\frac{3 k \lambda \sigma+k(k-1) \sigma^{2}}{6 \lambda+3 k \sigma}>0 \tag{5.30}
\end{equation*}
$$

For $0 \neq r \in \mathbb{Z}$ we denote by $\phi_{t}^{r}$ the diffeomorphism of $M$ composition

$$
\overbrace{\phi_{t} \circ \cdots \circ \phi_{t}}^{r}
$$

if $r>0$, and the obvious composition when $r<0$. By (5.22) one has $R_{\phi^{r}}=R_{\phi}$. From Theorem 10 together with (5.30) one deduces the following corollary.

Corollary 14. If $r, r^{\prime}$ are nonzero integers, then $\phi_{t}^{r}$ and $\tilde{\phi}_{t}^{r^{\prime}}$ are loops in $\operatorname{Ham}(M)$ which are not homotopic by a homotopy consisting of Hamiltonian circle actions.

Finally we consider the 1-parameter subgroup $\zeta$ in $\operatorname{Ham}(M)$ defined by the toric structure and the inclusion

$$
y \in S^{1} \mapsto\left(y^{l}, y^{\tilde{l}}\right) \in\left(S^{1}\right)^{2}
$$

where $l, \tilde{l} \in \mathbb{Z} \backslash\{0\}$. That is,

$$
\begin{equation*}
\zeta_{t}[z]=\left[z_{1} \mathrm{e}^{2 \pi \mathrm{i} l t}, z_{2} \mathrm{e}^{2 \pi \mathrm{i} \tilde{\mathrm{i}} t}, z_{3}, z_{4}\right] . \tag{5.31}
\end{equation*}
$$

The fixed point set $F$ of $\zeta$ is the singleton set $F=\left\{[z] \in M: z_{1}=z_{2}=0\right\}$. This point belongs to $U_{1}$; and in the coordinates $w_{1}, w_{1}^{\prime}$ (see (5.15)) on $U_{1}$

$$
\zeta_{t}\left(w_{1}, w_{1}^{\prime}\right)=\left(w_{1} \mathrm{e}^{2 \pi \mathrm{i} l t}, w_{1}^{\prime} \mathrm{e}^{2 \pi \mathrm{i} \tilde{I} t}\right)
$$

Hence

$$
\zeta_{t *}\left(\frac{\partial}{\partial w_{1}} \wedge \frac{\partial}{\partial w_{1}^{\prime}}\right)=\mathrm{e}^{2 \pi \mathrm{i}(l+\tilde{l}) t} \frac{\partial}{\partial w_{1}} \wedge \frac{\partial}{\partial w_{1}^{\prime}},
$$

and the corresponding function $L$ is the constant $2 \pi \mathrm{i}(l+\tilde{l})$.
On the other hand, the corresponding Hamiltonian vector field $Y_{M}$ is

$$
Y_{M}=-2 \pi \mathrm{i}\left(l w_{1} \frac{\partial}{\partial w_{1}}+\tilde{l} w_{1}^{\prime} \frac{\partial}{\partial w_{1}^{\prime}}\right),
$$

which vanishes on $F$. The normalized Hamiltonian function defined by $\zeta$ is

$$
l\left(\pi\left|z_{1}\right|^{2}-\kappa\right)+\tilde{l}\left(\pi\left|z_{2}\right|^{2}-\tilde{\kappa}\right)
$$

On $F$ this Hamiltonian reduces to the constant $-(l \kappa+\tilde{l} \tilde{\kappa})$.
According to our conventions the $U(1)$-action on $\mathbb{C} \frac{\partial}{\partial w_{1}}$ has multiplicity $-l$, and the multiplicity on the space $\mathbb{C} \frac{\partial}{\partial w_{1}^{\prime}}$ is $-\tilde{l}$, then the equivariant Euler class of the normal bundle to $F$ in $M$ is $e_{F}=l \tilde{l} u^{2}$. Thus by (5.1) and (5.3), $R_{\zeta}=(l+\tilde{l})(l \kappa+\tilde{l} \tilde{\kappa})(l \tilde{l})^{-1}$.

One can state the following proposition.
Proposition 15. Let $l$, $\tilde{l}$ be nonzero integers, and $\zeta$ the 1 -subgroup of $\operatorname{Ham}(M)$ defined by (5.31), then

$$
R_{\zeta}=\frac{(l+\tilde{l})(l \kappa+\tilde{l} \tilde{\kappa})}{l \tilde{l}}
$$

where the constants $\kappa$ and $\tilde{\kappa}$ are given by (5.19) and by (5.24) respectively.
$R_{\zeta}$ is a rational function in the variables $l, \tilde{l}$. Hence $R_{\zeta}=R_{\zeta^{r}}$, for any $r \in \mathbb{Z} \backslash\{0\}$. If $\mathbf{I}=(l, \tilde{l})$ and $\mathbf{I}^{\prime}=\left(l^{\prime}, \tilde{l}^{\prime}\right)$ are two pairs of nonzero integers such that the corresponding 1-parameter subgroups $\zeta(\mathbf{l})$ and $\zeta\left(\mathbf{I}^{\prime}\right)$ satisfy $R_{\zeta(\mathbf{I})} \neq R_{\zeta\left(\mathbf{I}^{\prime}\right)}$, then, by Theorem 10, $\zeta(\mathbf{l})^{r}$ and $\zeta\left(\mathbf{I}^{\prime}\right)^{s}$ are not homotopic by a family of Hamiltonian $S^{1}$-actions, whenever $r, s \in \mathbb{Z} \backslash\{0\}$.

When $M$ is the Hirzebruch surface determined by the triple ( $k=1, \tau>2, \sigma=1$ ), Abreu and McDuff proved in [2] that $\pi_{1}(\operatorname{Ham}(M))$ is isomorphic to $\mathbb{Z}$. Thus we have the following corollary.

Corollary 16. Let $M$ be the Hirzebruch surface defined by $(k=1, \tau>2, \sigma=1)$. There are infinitely many pairs $\left(\zeta(\mathbf{l}), \zeta\left(\mathbf{I}^{\prime}\right)\right)$ of 1-parameter closed subgroups of $\operatorname{Ham}(M)$ such that $[\zeta(\mathbf{l})]=\left[\zeta\left(\mathbf{I}^{\prime}\right)\right] \in \pi_{1}(\operatorname{Ham}(M))$, but

$$
[\zeta(\mathbf{l})] \neq\left[\zeta\left(\mathbf{I}^{\prime}\right)\right] \in[U(1), \operatorname{Ham}(M)]_{g h} .
$$

Remark 1. Two Hamiltonian circle actions on $M, \phi$ and $\phi^{\prime}$ are conjugate if there exists an element $h \in \operatorname{Ham}(M)$, such that $h \cdot \phi_{t} \cdot h^{-1}=\phi_{t}^{\prime}$ for all $t$. If $\phi$ and $\phi^{\prime}$ are conjugate, let $h_{s}$ be a path in $\operatorname{Ham}(M)$ from Id to $h$, then $h_{s} \cdot \phi \cdot h_{s}^{-1}$ defines a homotopy between $\phi$ and $\phi^{\prime}$, and $[\phi]=\left[\phi^{\prime}\right] \in[U(1), \operatorname{Ham}(M)]_{g h}$. By Corollary 16, there are infinitely many conjugacy classes of circle actions on the Hirzebruch surface considered in this corollary.

Remark 2. Although the characteristic $R_{1}$ allows us to distinguish infinitely many conjugacy classes of Hamiltonian circle actions in a Hirzebruch surface $M$, the situation is different for $U(1)^{2}$-actions, as we show next.

Given $\mathbf{I}=(l, \tilde{l})$ a pair of nonzero integers, we define a $U(1)^{2}$-action on $M$ by

$$
\xi_{s t}[z]=\left[z_{1} \mathrm{e}^{2 \pi \mathrm{i} l s}, z_{2} \mathrm{e}^{2 \pi \mathrm{i} \tilde{l} t}, z_{3}, z_{4}\right] .
$$

The fixed point set is again the singleton $F$, and the equivariant Euler class $e_{F}=l \tilde{l} u v$, where $u, v$ are coordinates on $\mathfrak{u}(1) \oplus \mathfrak{u}(1)$. According to the localization formula (5.1), $R_{1}(\xi)$ is a rational function of the variables $\check{u}:=l u, \check{v}:=\tilde{l} v$. The numerator is a homogeneous degree two polynomial $A_{1} \check{u}^{2}+A_{2} \check{u} \check{v}+A_{3} \check{v}^{2}$, with $A_{j}$ independent of $\mathbf{l}$; and the denominator is $\check{u} \check{v}$. As $R_{1}(\xi) \in H^{0}\left(B\left(U(1)^{2}\right)\right)$, then $A_{1}=A_{3}=0$, and $R_{1}(\xi)$ is independent of $\mathbf{l}$. That is, $R_{1}$ is constant on $\{\xi(\mathbf{I})\}_{1}$. But Karshon proved that the number of conjugacy classes of maximal tori in $M$ is the smallest integer greater than or equal to $\frac{\lambda}{\sigma}$ (see [15]). That is, the characteristic number $R_{1}$ is not fine enough to analyze this case.

## Acknowledgements

This work has been partially supported by Ministerio de Ciencia y Tecnología, grant MAT2003-09243-C02-00.
I thank Dusa McDuff for her enlightening comments. I thank an anonymous referee for constructive comments and for having pointed out for me the references [18,15]. Remark 1 to Corollary 16 is really a comment of the referee.

## References

[1] M. Abreu, Topology of symplectomorphism group of $S^{2} \times S^{2}$, Invent. Math. 131 (1998) 1-23.
[2] M. Abreu, D. McDuff, Topology of symplectomorphism groups of rational ruled surfaces, J. Amer. Math. Soc. 13 (2000) $971-1009$.
[3] M.F. Atiyah, R. Bott, The moment map and equivariant cohomology, Topology 23 (1984) 1-28.
[4] M. Audin, Torus Actions on Symplectic Manifolds, Birkhäuser, Basel, 2004.
[5] N. Berline, E. Getzler, M. Vergne, Heat Kernels and Dirac Operators, Springer-Verlag, Berlin, 1991.
[6] C.J. Earle, J. Eells, A fibre bundle description of Teichmüller theory, J. Differential Geom. 3 (1969) 19-43.
[7] W. Fulton, J. Harris, Representation Theory, Springer-Verlag, New York, 1991.
[8] R. Goodman, N.R. Wallach, Representations and Invariants of Classical Groups, Cambridge U.P., Cambridge, 1998.
[9] M. Gromov, Pseudo holomorphic curves in symplectic manifolds, Invent. Math. 82 (1985) 307-347.
[10] V. Guillemin, Moment Maps and Combinatorial Invariants of Hamiltonian $T^{n}$-Spaces, Birkhäuser, Boston, 1994.
[11] V. Guillemin, L. Lerman, S. Sternberg, Symplectic Fibrations and Multiplicity Diagrams, Cambridge U.P., Cambridge, 1996.
[12] V. Guillemin, S. Sternberg, Symplectic Techniques in Physics, Cambridge U.P., Cambridge, 1984.
[13] V. Guillemin, S. Sternberg, Supersymmetry and Equivariant de Rham Theory, Springer-Verlag, Berlin, 1999.
[14] T. Januszkiewicz, J. Kedra, Characteristic classes of smooth fibrations. ArXiv: math.SG/0209288.
[15] Y. Karshon, Maximal tori in the symplectomorphism groups of Hirzebruch surfaces, Math. Res. Lett. 10 (2003) $125-132$.
[16] J. Kedra, D. McDuff, Homotopy properties of Hamiltonian group actions, Geom. Topol. 9 (2005) 121-162.
[17] A.A. Kirilov, Elements of the Theory of Representations, Springer-Verlag, Berlin, 1976.
[18] F. Lalonde, D. McDuff, L. Polterovich, On the flux conjectures, in: CRM Proceedings and Lecture Notes 15, Amer. Math. Soc., Providence, RI, 1998, pp. 69-85.
[19] F. Lalonde, D. McDuff, L. Polterovich, Topological rigidity of Hamiltonian loops and quantum homology, Invent. Math. 135 (1999) $369-385$.
[20] D. McDuff, Lectures on groups of symplectomorphisms. ArXiv: math.SG/0201032 (preprint).
[21] D. McDuff, D. Salamon, Introduction to Symplectic Topology, Clarenton Press, Oxford, 1998.
[22] D. McDuff, D. Salamon, J-Holomorphic Curves and Symplectic Topology, Amer. Math. Soc. Colloq. Publ., Providence, 2004.
[23] D. McDuff, S. Tolman, On nearly semifree circle actions. ArXiv: math.SG/0503467 (preprint).
[24] D. McDuff, S. Tolman, Polytopes with mass linear functions (in preparation).
[25] L. Polterovich, The Geometry of the Group of Symplectic Diffeomorphisms, Birkhäuser, Basel, 2001.
[26] D. Salamon, E. Zehnder, Morse theory for periodic solutions of Hamiltonian systems and the Maslov index, Comm. Pure Appl. Math. XLV (1992) 1303-1360.
[27] A. Viña, Generalized symplectic action and symplectomorphism groups of coadjoint orbits, Ann. Global Anal. Geom. 28 (2005) $309-318$.
[28] A. Viña, A characteristic number of Hamiltonian bundles over $S^{2}$, J. Geom. Phys. 56 (2006) 2327-2343.
[29] A. Weinstein, Cohomology of symplectomorphism groups and critical values of Hamiltonians, Math. Z. 201 (1989) 75-82.


[^0]:    E-mail address: vina@uniovi.es.

